

# 耦合 Klein-Gordon-Schrödinger 方程新的精确解

刘成仕<sup>†</sup> 杜兴华

(大庆石油学院数学系, 大庆 163318)

(2004 年 7 月 9 日收到, 2004 年 11 月 26 日收到修改稿)

用函数变换法将耦合 Klein-Gordon-Schrödinger 方程化成可求解的不定积分形式, 进而求出其精确解, 包括有理函数型解、三角函数型解、双曲函数型解, 以及椭圆函数型解. 所得结果包含了文献中用齐次平衡法所得的解为特例.

关键词: 耦合 Klein-Gordon-Schrödinger 方程, 精确解, 椭圆函数, 双曲函数

PACC: 0340K, 0290

## 1. 引 言

耦合 Klein-Gordon-Schrödinger 方程(以下简称 cKGS 方程)

$$i \frac{\partial}{\partial t} \psi + \frac{1}{2}(\phi_{xx} + \phi_{yy} + \phi_{zz}) + \phi\psi = 0, \quad (1)$$

$$\frac{\partial^2}{\partial t^2} \phi - (\phi_{xx} + \phi_{yy} + \phi_{zz}) + m^2 \phi - |\psi|^2 = 0 \quad (2)$$

描述了一个保守的复原子核场  $\psi$  与一个实中性介子场  $\phi$  的相互作用的古典 Yukawa 型.  $\psi$  和  $\phi$  均为标量场,  $m$  为介子的质量<sup>[1,2]</sup>. Zhang<sup>[3]</sup>用齐次平衡法得到了方程(1)和(2)的显解

$$\psi(x, y, z, t) = \frac{3}{4} k^2 \operatorname{sech}^2\left(\frac{kx + ly + pz \pm mt}{2}\right), \quad (3)$$

$$\begin{aligned} \phi(x, y, z, t) = & \frac{3}{4} k^2 \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2} \delta t\right)\right] \\ & \times \operatorname{sech}^2\left(\frac{kx + ly + pz \pm mt}{2}\right), \quad (4) \end{aligned}$$

其中

$$k^2 = \frac{6}{a} m^2, \quad \alpha_1 k + \beta_1 l + \gamma_1 p = \pm m,$$

$$\alpha_1^2 + \beta_1^2 + \gamma_1^2 + \delta = 2m^2, \quad k^2 + l^2 + p^2 = 2m^2, \quad (5)$$

$a$  是任一大于零的常数, 这有误. 将(3)和(4)式代入(1)和(2)式, 利用(5)式即得  $a=3$

本文利用函数变换法化方程(1)和(2)为可解的不定积分形式, 求出其精确解, 包括有理函数型解、三角函数型解、双曲函数型解和椭圆函数型解, 并以

(3)和(4)式作为特例.

## 2. cKGS 方程(1)和(2)的精确解

为求方程(1)和(2)的精确解, 如文献[3], 假设

$$\psi(x, y, z, t) = \psi(x, y, z, t) \exp[i g(x, y, z, t)], \quad (6)$$

$$g(x, y, z, t) = \alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2} \delta t, \quad (7)$$

则方程(1)和(2)成为

$$\psi_u - (\phi_{xx} + \phi_{yy} + \phi_{zz}) + m^2 \phi - \psi^2 = 0, \quad (8)$$

$$\begin{aligned} & \frac{1}{2}(\phi_{xx} + \phi_{yy} + \phi_{zz}) \\ & - \frac{1}{2}(\alpha_1^2 + \beta_1^2 + \gamma_1^2 + \frac{1}{2} \delta) \psi + \psi^2 = 0, \quad (9) \end{aligned}$$

$$\psi_l + \alpha_1 \psi_x + \beta_1 \psi_y + \gamma_1 \psi_z = 0. \quad (10)$$

设(8)–(10)式的行波解为

$$\psi(x, y, z, t) = \psi(\xi), \quad (11)$$

其中

$$\xi = kx + ly + pz + \omega t, \quad (12)$$

下面求方程(1)和(2)的形如(6)–(7)–(11)和(12)式的精确解. 将(11)式代入(8)–(10)式, 有

$$(\omega^2 - (k^2 + l^2 + p^2)) \psi_{\xi\xi} + m^2 \psi - \psi^2 = 0, \quad (13)$$

$$\begin{aligned} & \frac{1}{2}(k^2 + l^2 + p^2) \psi_{\xi\xi} \\ & - \frac{1}{2}(\alpha_1^2 + \beta_1^2 + \gamma_1^2 + \delta) \psi + \psi^2 = 0, \quad (14) \end{aligned}$$

$$(\omega + \alpha_1 k + \beta_1 l + \gamma_1 p) \psi_{\xi} = 0. \quad (15)$$

<sup>†</sup>E-mail: chengshiliu-68@126.com

由 (13) 和 (14) 式相对照, 再令 (15) 式中  $\phi_\xi \neq 0$  知若

$$k^2 + l^2 + p^2 = 2\omega^2, \quad (16)$$

$$\frac{1}{2}(\alpha_1^2 + \beta_1^2 + \gamma_1^2 + \delta) = m^2, \quad (17)$$

$$\omega + \alpha_1 k + \beta_1 l + \gamma_1 p = 0 \quad (18)$$

成立, 则只需考虑:

$$\phi_{\xi\xi} = -\frac{1}{\omega^2}\phi^2 + \frac{m^2}{\omega^2}\phi. \quad (19)$$

对 (19) 式积分一次, 有

$$\phi_\xi^2 = -\frac{2}{3\omega^2}\phi^3 + \frac{m^2}{\omega^2}\phi^2 + c, \quad (20)$$

$c$  为积分常数. 令

$$W = \left(-\frac{2}{3\omega^2}\right)^{1/3}\phi, \quad \xi_1 = \left(-\frac{2}{3\omega^2}\right)^{1/3}\xi, \quad (21)$$

则 (20) 式写成

$$\left(\frac{dW}{d\xi_1}\right)^2 = F(W) = W^3 + aW^2 + c, \quad (22)$$

其中

$$a = \frac{m^2}{\omega^2}\left(-\frac{2}{3\omega^2}\right)^{-2/3}. \quad (23)$$

下面再分两种情形讨论:

情形 1  $\Delta = 0$  这里  $\Delta = -27\left(\frac{2}{27}\alpha^3 + c\right)^2 - 4 \times$

$\left(-\frac{a^2}{3}\right)^3$  为  $F(W)$  的判别式. 此时  $F(W) = 0$  有重根, 设为  $F(W) = (W - \alpha)(W - \beta)$ , 由 (22) 式知, 当  $W > \beta$  时, 有以下三种情形:

$$\begin{aligned} \pm(\xi_1 - \xi_0) &= \int \frac{dW}{(W - \alpha)\sqrt{W - \beta}} \\ &= \frac{1}{\sqrt{\alpha - \beta}} \ln \left| \frac{\sqrt{W - \beta} - \sqrt{\alpha - \beta}}{\sqrt{W - \beta} + \sqrt{\alpha - \beta}} \right|, \quad \alpha > \beta, \end{aligned} \quad (24)$$

$$\pm(\xi_1 - \xi_0) = \frac{2}{\sqrt{\beta - \alpha}} \arctan \sqrt{\frac{W - \beta}{\beta - \alpha}}, \quad \alpha < \beta, \quad (25)$$

$$\pm(\xi_1 - \xi_0) = -\frac{2}{\sqrt{W - \beta}}, \quad \alpha = \beta, \quad (26)$$

其中  $\xi_0$  为积分常数. 相应地 (19) 式的解为

$$\begin{aligned} \phi_1 &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \\ &\times \left[ \left( \frac{2\sqrt{\alpha - \beta}}{1 + \exp\left[\pm\sqrt{\alpha - \beta}\left(-\frac{2}{3\omega^2}\right)^{1/3}(\xi - \xi_0)\right]} \right) \right. \end{aligned}$$

$$\begin{aligned} &\left. - \sqrt{\alpha - \beta} \right)^2 + \beta \Big] = \left(-\frac{2}{3\omega^2}\right)^{-1/3} \\ &\times \left[ (\alpha - \beta) \operatorname{th}^2\left(\frac{1}{2}\sqrt{\alpha - \beta}\left(-\frac{2}{3\omega^2}\right)^{1/3}\right. \right. \\ &\left. \left. \times (\xi - \xi_0)\right) + \beta \right], \quad \alpha > \beta, \quad (27) \end{aligned}$$

$$\begin{aligned} \phi_2 &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \\ &\times \left[ \left( \frac{2\sqrt{\alpha - \beta}}{1 - \exp\left[\pm\sqrt{\alpha - \beta}\left(-\frac{2}{3\omega^2}\right)^{1/3}(\xi - \xi_0)\right]} \right) \right. \\ &\left. - \sqrt{\alpha - \beta} \right)^2 + \beta \Big] \\ &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \left[ (\alpha - \beta) \operatorname{cth}^2\left(\frac{1}{2}\sqrt{\alpha - \beta}\right. \right. \\ &\left. \left. \times \left(-\frac{2}{3\omega^2}\right)^{1/3}(\xi - \xi_0)\right) + \beta \right], \quad \alpha > \beta, \quad (28) \end{aligned}$$

$$\begin{aligned} \phi_3 &= \left(-\frac{2}{3\omega^2}\right)^{-1/3}(\beta - \alpha) \\ &\times \sec^2\left(\frac{\sqrt{\beta - \alpha}}{2}\left(-\frac{2}{3\omega^2}\right)^{1/3}(\xi - \xi_0)\right) \\ &+ \left(-\frac{2}{3\omega^2}\right)^{1/3}\alpha, \quad \alpha < \beta, \quad (29) \end{aligned}$$

$$\begin{aligned} \phi_4 &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \left[ \beta + \frac{4}{\left(-\frac{2}{3\omega^2}\right)^{2/3}(\xi - \xi_0)^2} \right], \\ &\alpha = \beta. \quad (30) \end{aligned}$$

(27)–(30) 式中积分常数  $\xi_0$  已经重新标度, 仍用  $\xi_0$  表示, 下同. 这里 (27) 和 (28) 式是 (19) 式的孤波解, (29) 式是三角函数型周期解, (30) 式是其有理函数解. 相应地可以写出方程 (1) 和 (2) 的显式精确解

$$\begin{aligned} &\phi_1(x, y, z, t) \\ &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \left[ (\alpha - \beta) \operatorname{th}^2\left(\frac{1}{2}\sqrt{\alpha - \beta}\left(-\frac{2}{3\omega^2}\right)^{1/3}\right. \right. \\ &\left. \left. \times (\xi - \xi_0)\right) + \beta \right], \quad \alpha > \beta, \quad (31) \end{aligned}$$

$$\begin{aligned} &\psi_1(x, y, z, t) \\ &= \phi_1(x, y, z, t) \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2}\delta t\right)\right], \end{aligned} \quad (32)$$

$$\phi_2(x, y, z, t)$$

$$= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \left[ (\alpha - \beta) \operatorname{th}^2\left(\frac{1}{2} \sqrt{\alpha - \beta} \left(-\frac{2}{3\omega^2}\right)^{1/3} \times (\xi - \xi_0)\right) + \beta \right], \quad \alpha > \beta, \quad (33)$$

$$\psi_2(x, y, z, t) = \phi_2(x, y, z, t) \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2} \delta t\right)\right], \quad (34)$$

$$\phi_3(x, y, z, t) = \left(-\frac{2}{3\omega^2}\right)^{-1/3} (\beta - \alpha) \operatorname{sec}^2\left[\frac{\sqrt{\beta - \alpha}}{2} \left(-\frac{2}{3\omega^2}\right)^{1/3} \times (\xi - \xi_0)\right] + \left(-\frac{2}{3\omega^2}\right)^{1/3} \alpha, \quad \alpha < \beta, \quad (35)$$

$$\psi_3(x, y, z, t) = \phi_3(x, y, z, t) \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2} \delta t\right)\right], \quad (36)$$

$$\phi_4 = \left(-\frac{2}{3\omega^2}\right)^{-1/3} \times \left(\beta + \frac{4}{\left(-\frac{2}{3\omega^2}\right)^{-2/3} (\xi - \xi_0)}\right), \quad \alpha = \beta, \quad (37)$$

$$\psi_4(x, y, z, t) = \phi_4(x, y, z, t) \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2} \delta t\right)\right]. \quad (38)$$

注 1 事实上 (24) 和 (25) 式对应于  $a \neq 0$  (26) 式对应于  $a = 0$ , 此时  $F(W) = W^3$ , 即  $\alpha = \beta = 0$ .

注 2 当  $c = 0$  时, 有  $\Delta = 0$ . 此时有  $\alpha = 0, \beta = -\frac{m^2}{\omega^2} \left(-\frac{2}{3\omega^2}\right)^{-2/3}$ . 若实数  $m \neq 0$  则  $\alpha \neq \beta$ . 由 (27) — (29) 式, 有

$$\phi(\xi) = \frac{3}{2} m^2 \left[ 1 - \operatorname{th}^2\left(\frac{m}{\omega} (\xi - \xi_0)\right) \right], \quad \omega^2 > 0, \quad (39)$$

$$\phi(\xi) = \frac{3}{2} m^2 \left[ 1 - \operatorname{cth}^2\left(\frac{m}{\omega} (\xi - \xi_0)\right) \right], \quad \omega^2 > 0, \quad (40)$$

$$\phi(\xi) = \frac{3}{2} m^2 \operatorname{sec}^2\left(\frac{1}{2} \sqrt{-\frac{m^2}{\omega^2}} (\xi - \xi_0)\right), \quad \omega^2 < 0. \quad (41)$$

在 (39) 式中, 令  $\omega = m, k^2 = 2m^2, \xi_0 = 0$ , 有

$$\phi(\xi) = \frac{3}{4} k^2 \operatorname{sech}^2\left(\frac{1}{2}(kx + ly + pz \pm \omega t)\right). \quad (42)$$

此式即为 (3) 式. 因此文献 [3] 中的结果只是 (31) 和

(32) 式中的参数取特殊值时的特例.

情形 2  $\Delta \neq 0$ , 再分两种情形讨论.

1)  $\Delta > 0, a \neq 0$ . 此时  $F(W) = 0$  有三个不同的实根  $\alpha < \beta < \gamma$ , 当  $\alpha < W < \beta$  时, 作变量替换

$$W = \alpha + (\beta - \alpha) \sin^2 \varphi. \quad (43)$$

由 (22) 式, 有

$$\begin{aligned} & \pm (\xi_1 - \xi_0) \\ &= \int \frac{dW}{\sqrt{F(W)}} \\ &= \int \frac{\alpha(\beta - \alpha) \sin \varphi \cos \varphi d\varphi}{\sqrt{\gamma - \alpha} (\beta - \alpha) \sin \varphi \cos \varphi \sqrt{1 - k^2 \sin^2 \varphi}} \\ &= \frac{2}{\sqrt{\gamma - \alpha}} \int \frac{d\varphi}{\sqrt{1 - k^2 \sin^2 \varphi}}, \quad (44) \end{aligned}$$

其中  $k^2 = \frac{\beta - \alpha}{\gamma - \alpha}$ . 由 (44) 式, 有

$$\pm \frac{\sqrt{\gamma - \alpha}}{2} (\xi_1 - \xi_0) = \int \frac{d\varphi}{\sqrt{1 - k^2 \sin^2 \varphi}}. \quad (45)$$

由 (43) 和 (45) 式及 Jacobi 椭圆正弦函数的定义<sup>[5]</sup> 知,

$$\operatorname{sn}^2\left(\pm \frac{\sqrt{\gamma - \alpha}}{2} (\xi_1 - \xi_0), k\right) = \sin^2 \varphi = \frac{W - \alpha}{\beta - \alpha}, \quad (46)$$

即

$$W = \alpha + (\beta - \alpha) \operatorname{sn}^2\left(\frac{\sqrt{\gamma - \alpha}}{2} (\xi_1 - \xi_0), k\right), \quad (47)$$

相应地方程 (19) 的椭圆周期解为

$$\begin{aligned} \phi_5(\xi) &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \alpha + \left(-\frac{2}{3\omega^2}\right)^{-1/3} (\beta - \alpha) \\ &\times \operatorname{sn}^2\left(\frac{\sqrt{\gamma - \alpha}}{2} \left(-\frac{2}{3\omega^2}\right)^{-1/3} (\xi - \xi_0), k\right). \quad (48) \end{aligned}$$

由此得 cKGS 方程 (1) 和 (2) 的含椭圆函数的精确解为

$$\begin{aligned} \phi_5(\xi) &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \alpha + \left(-\frac{2}{3\omega^2}\right)^{-1/3} (\beta - \alpha) \\ &\times \operatorname{sn}^2\left(\frac{\sqrt{\gamma - \alpha}}{2} \left(-\frac{2}{3\omega^2}\right)^{-1/3} (\xi - \xi_0), k\right), \quad (49) \end{aligned}$$

$$\psi_5(x, y, z, t)$$

$$= \phi_5(x, y, z, t) \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2} \delta t\right)\right]. \quad (50)$$

若  $W > \gamma$ , 则作变量代换

$$W = \frac{-\beta \sin^2 \phi + \gamma}{\cos^2 \phi}, \quad (51)$$

相似地,方程(19)的解为

$$\begin{aligned} & \phi_6(\xi) \\ &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \\ & \times \frac{-\beta \operatorname{sn}^2\left(\frac{\sqrt{\gamma-\alpha}}{2}\left(-\frac{2}{3\omega^2}\right)^{-1/3}(\xi-\xi_0), k\right) + \gamma}{\operatorname{cn}^2\left(\frac{\sqrt{\gamma-\alpha}}{2}\left(-\frac{2}{3\omega^2}\right)^{-1/3}(\xi-\xi_0), k\right)}, \end{aligned} \quad (52)$$

其中  $k^2 = \frac{\beta-\alpha}{\gamma-\alpha}$ , 由此得 cKGS 方程(1)和(2)的含椭圆函数的精确解为

$$\begin{aligned} & \phi_6(\xi) \\ &= \left(-\frac{2}{3\omega^2}\right)^{-1/3} \\ & \times \frac{-\beta \operatorname{sn}^2\left(\frac{\sqrt{\gamma-\alpha}}{2}\left(-\frac{2}{3\omega^2}\right)^{-1/3}(\xi-\xi_0), k\right) + \gamma}{\operatorname{cn}^2\left(\frac{\sqrt{\gamma-\alpha}}{2}\left(-\frac{2}{3\omega^2}\right)^{-1/3}(\xi-\xi_0), k\right)}, \end{aligned} \quad (53)$$

$$\begin{aligned} & \psi_6(x, y, z, t) \\ &= \phi_6(x, y, z, t) \\ & \times \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2}\delta t\right)\right]. \end{aligned} \quad (54)$$

2)  $\Delta < 0$ , 此时  $F(W) = 0$  只有一个实根, 设为  $F(W) = (W - \alpha) \sqrt{W^2 + p_0 W + q_0}$ , 且  $p_0^2 - 4q_0 < 0$ , 当  $W > \alpha$  时, 作变量替换

$$W = \alpha + \sqrt{\alpha^2 + p_0 \alpha + q_0} \tan^2 \frac{\varphi}{2}. \quad (55)$$

由(22)式, 有

$$\begin{aligned} & \pm(\xi_1 - \xi_0) \\ &= \int \frac{dW}{\sqrt{(W - \alpha) \sqrt{W^2 + p_0 W + q_0}}} \\ &= \int \frac{\sqrt{\alpha^2 + p_0 \alpha + q_0} \tan \frac{\varphi}{2}}{\cos^2 \frac{\varphi}{2}} d\varphi \\ &= \int \frac{\tan \frac{\varphi}{2}}{(\alpha^2 + p_0 \alpha + q_0)^{1/4} \sqrt{1 - k^2 \sin^2 \varphi}} \\ &= \frac{1}{(\alpha^2 + p_0 \alpha + q_0)^{1/4}} \int \frac{d\varphi}{\sqrt{1 - k^2 \sin^2 \varphi}}, \end{aligned} \quad (56)$$

其中  $k^2 = \frac{1}{2} \left(1 - \frac{\alpha + \frac{p_0}{2}}{\sqrt{\alpha^2 + p_0 \alpha + q_0}}\right)$ . 由(56)式及

Jacobi 椭圆余弦函数的定义<sup>[5]</sup>知,

$$\operatorname{cn}\left((\alpha^2 + p_0 \alpha + q_0)^{1/4}(\xi_1 - \xi_0), k\right) = \cos \varphi. \quad (57)$$

又由变量代换(55)式, 有

$$\cos \varphi = \frac{2\sqrt{\alpha^2 + p_0 \alpha + q_0}}{W - \alpha + \sqrt{\alpha^2 + p_0 \alpha + q_0}} - 1. \quad (58)$$

由(57)和(58)式知, 当  $W > \alpha$  时, 有

$$\begin{aligned} W &= \alpha + \frac{2\sqrt{\alpha^2 + p_0 \alpha + q_0}}{1 + \operatorname{cn}\left((\alpha^2 + p_0 \alpha + q_0)^{1/4}(\xi_1 - \xi_0), k\right)} \\ & \quad - \sqrt{\alpha^2 + p_0 \alpha + q_0}. \end{aligned} \quad (59)$$

相应地, 方程(19)的椭圆周期解为

$$\phi_7(\xi) = \left(-\frac{2}{3\omega^2}\right)^{-1/3} \left[ \alpha + \frac{2\sqrt{\alpha^2 + p_0 \alpha + q_0}}{1 + \operatorname{cn}\left((\alpha^2 + p_0 \alpha + q_0)^{1/4}\left(-\frac{2}{3\omega^2}\right)^{1/3}(\xi - \xi_0), k\right)} - \sqrt{\alpha^2 + p_0 \alpha + q_0} \right]. \quad (60)$$

由(60)式可写出 cKGS 方程(1)和(2)的另一含椭圆函数的精确解

$$\phi_7(\xi) = \left(-\frac{2}{3\omega^2}\right)^{-1/3} \left[ \alpha + \frac{2\sqrt{\alpha^2 + p_0 \alpha + q_0}}{1 + \operatorname{cn}\left((\alpha^2 + p_0 \alpha + q_0)^{1/4}\left(-\frac{2}{3\omega^2}\right)^{1/3}(\xi - \xi_0), k\right)} - \sqrt{\alpha^2 + p_0 \alpha + q_0} \right], \quad (61)$$

$$\psi_7(x, y, z, t) = \phi_7(x, y, z, t) \exp\left[i\left(\alpha_1 x + \beta_1 y + \gamma_1 z + \frac{1}{2}\delta t\right)\right]. \quad (62)$$

### 3. 结束语

尽管已经建立了许多求解非线性方程的方法,

如齐次平衡法<sup>[4, 6, 7]</sup>、试探函数法<sup>[8, 9]</sup>、双曲函数展开法<sup>[10-12]</sup>、椭圆函数展开法<sup>[13-15]</sup>、函数变换法<sup>[16]</sup>、sine-cosine 法<sup>[17-19]</sup>等, 这一领域依然充满挑战. 本文的方法得到的结果是系统的, 得到了在条件(16)一

(18) 式下, 形如 (6) (7) (11) 和 (12) 式的全部精确解  $\phi_i, \psi_i (i = 1, \dots, 7)$ .

- [ 1 ] Tsutsumi F I 1978 *J. Math. Anal. Appl.* **66** 358
- [ 2 ] Yukawa H 1935 *Proc. Phys. -Math. Soc. Japan* **17** 48
- [ 3 ] Zhang H Q 2002 *Acta Math. Sci.* A **22** 332 [ in Chinese ] 张辉群 2002 数学物理学报 A **22** 332 ]
- [ 4 ] Wang M L, Zhou Y B and Li Z B 1996 *Phys. Lett. A* **216** 67
- [ 5 ] Wang Z X and Guo D R 2002 *Special Functions* ( Beijing :Peking University Press ) [ in Chinese ] 王竹溪、郭敦仁 2002 特殊函数概论 (北京 北京大学出版社 )
- [ 6 ] Fan E G and Zhang H Q 1998 *Acta Phys. Sin.* **47** 353 [ in Chinese ] [ 范恩贵、张鸿庆 1998 物理学报 **47** 353 ]
- [ 7 ] Zhang J F 1999 *Chin. Phys. Lett.* **16** 4
- [ 8 ] Kudryashov N A 1990 *Phys. Lett. A* **147** 287
- [ 9 ] Zhang W G 2003 *Acta Math. Sci.* A **23** 679 [ in Chinese ] 张卫国 2003 数学物理学报 A **23** 679 ]
- [ 10 ] Li Z B and Zhang S Q 1997 *Acta Math. Sin.* **17** 81 [ in Chinese ] [ 李志斌、张善卿 1997 数学物理学报 **17** 81 ]
- [ 11 ] Lü K P, Shi Y R, Duan W S and Zhao J B 2001 *Acta Phys. Sin.* **50** 2074 [ in Chinese ] 吕克璞、石玉仁、段文山、赵金保 2001 物理学报 **50** 2074 ]
- [ 12 ] Zhang G X, Li Z B and Duan Y S 2000 *Sci. China A* **30** 1103 [ in Chinese ] 张桂茂、李志斌、段一士 2000 中国科学 A **30** 1103 ]
- [ 13 ] Liu S K, Fu Z T, Liu S D and Zhao Q 2001 *Acta Phys. Sin.* **50** 2068 [ in Chinese ] 刘式适、傅遵涛、刘式达、赵 强 2001 物理学报 **50** 2068 ]
- [ 14 ] Liu S K, Fu Z T, Liu S D and Zhao Q 2002 *Acta Phys. Sin.* **51** 718 [ in Chinese ] 刘式适、傅遵涛、刘式达、赵 强 2002 物理学报 **51** 718 ]
- [ 15 ] Porubov A V 1996 *Phys. Lett. A* **221** 391
- [ 16 ] Liu S K, Fu Z T, Liu S D and Zhao Q 2001 *Appl. Math. Mech.* **22** 326
- [ 17 ] Yan C T 1996 *Phys. Lett. A* **224** 77
- [ 18 ] Yan Z Y, Zhang H Q and Fan E G 1999 *Acta Phys. Sin.* **48** 1 [ in Chinese ] 闫振亚、张鸿庆、范恩贵 1999 物理学报 **48** 1 ]
- [ 19 ] Yan Z Y and Zhang H Q 1999 *Acta Phys. Sin.* **48** 1962 [ in Chinese ] 闫振亚、张鸿庆 1999 物理学报 **48** 1962 ]

## New exact solutions of coupled Klein-Gordon-Schrödinger equations

Liu Cheng-Shi Du Xing-Hua

( Department of Mathematics , Daqing Petroleum Institute , Daqing 163318 , China )

( Received 9 July 2004 ; revised manuscript received 26 November 2004 )

### Abstract

By using a function-transformation method , coupled Klein-Gordon-Schrödinger equations are reduced to the integrals that can be solved with elementary functions and elliptic functions. Its exact solutions that include discontinuous solutions , hyperbolic function type solutions and elliptic function type solutions , and a special solutions obtained in a reference , are given.

**Keywords :** coupled Klein-Gordon-Schrödinger equations , exact solutions , elliptic function , hyperbolic function

**PACC :** 0340K , 0290