

变系数广义 KdV 方程新的类孤波解和精确解

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用普通 KdV 方程作变换, 构造变系数广义 KdV 方程的解, 获得了变系数广义 KdV 方程新的 Jacobi 椭圆函数精确解、类孤波解、三角函数解和 Weierstrass 椭圆函数解.

关键词: KdV 方程, 变系数广义 KdV 方程, 类孤波解, 精确解

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1. 引 言

含外力的变系数广义 Korteweg-de Vries (KdV) 方程^[1]

$$u_{x,t} + 6F_1(u_x^2 + uu_{xx}) + F_2u_{xxx} + 6F_3u_x + F_4 = 0, \quad (1)$$

式中 $u_{x,t} = \frac{\partial^2 u}{\partial x \partial t}$, $u_{xx} = \frac{\partial^2 u}{\partial x^2}$, \dots , F_1, F_2, F_3, F_4 为 t 的函数. 如取(1)式中的外力项 $F_4 = 0$, 则(1)式演化为

$$u_{x,t} + 6F_1(u_x^2 + uu_{xx}) + F_2u_{xxx} + 6F_3u_x = 0. \quad (2)$$

方程(1)和(2)在流体力学、等离子体物理、气体动力学等领域有重要的应用^[1], 是物理学家和数学家感兴趣的方程之一. 为了求解非线性偏微分方程, 文献[1]利用 Bäcklund 变换, 得到了(1)式具有 Painlevé 性质条件下的解; 文献[2]用截断展开方法获得广义变系数方程新的精确类孤波解; 文献[3]用 Jacobi 椭圆函数展开法, 获得不含外力项的变系数 KdV 方程的解; 文献[4]利用 Bäcklund 变换研究了广义 KP 方程的类孤波解; 文献[5, 6]利用扩展的双曲正切函数法、文献[7]用反散射法研究了变系数存在外力项的 KdV 方程的精确解; 文献[8]用新的综合的扩展方法, 找到了变系数 KdV 方程的解; 最近, 文献[9]用扩展的 Jacobi 椭圆函数展开法, 获得了普通 KdV 方程新的解; 文献[10]获得了 KP 方程的新解; 文献[11, 12]用其他方法研究了非线性方程的解.

本文用普通 KdV 方程作变换, 构造含外力的变系数广义 KdV 方程的解, 获得了(1)和(2)式新的精确解、类孤波解、三角函数解和 Weierstrass 椭圆函数解.

2. 解的构造

对于普通的 KdV 方程

$$v_t + vv_x + \beta v_{xxx} = 0, \quad (3)$$

式中 β 为常数. 对 v 作行波变换 $v = v(\xi)$, $\xi = x - ct$, 代入(3)式, 并积分两次, 得^[13]

$$\frac{dv}{d\xi} = \frac{1}{3} \sqrt{-\frac{3v(\xi)^3 - 9cv(\xi)^2 - 18Av(\xi) - 18B}{\beta}}, \quad (4)$$

式中 c 为常数, A, B 为积分常数, $v(\xi)^2 \equiv v^2(\xi) \equiv (v(\xi))^2$. 当 A, B, c, β 取不同的值时, (4)式的 v 有下列的解:

1) 当 $B = 0$ (4)式的 v 有 9 组 Jacobi 椭圆函数解:

$$v_1 = \operatorname{sn}(\xi), A = -\frac{1}{6m^2}, c = \frac{1+m^2}{3m^2}, \beta = -\frac{1}{12m^2}, \quad (5)$$

$$v_2 = \operatorname{cn}(\xi), A = \frac{1-m^2}{6m^2}, c = \frac{2m^2-1}{3m^2}, \beta = \frac{1}{12m^2}, \quad (6)$$

$$v_3 = \operatorname{dn}(\xi), A = \frac{-1+m^2}{6},$$

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$$c = \frac{2 - m^2}{3}, \beta = \frac{1}{12}, \quad (7)$$

$$v_4 = a \operatorname{sn}(\xi)^2, A = -\frac{a^2}{6m^2},$$

$$c = \frac{(1 + m^2)a}{3m^2}, \beta = -\frac{a}{12m^2}, \quad (8)$$

$$v_5 = a \operatorname{cn}(\xi)^2, A = \frac{a^2(1 - m^2)}{6m^2},$$

$$c = \frac{(2m^2 - 1)a}{3m^2}, \beta = \frac{a}{12m^2}, \quad (9)$$

$$v_6 = a \operatorname{dn}(\xi)^2, A = \frac{a^2(-1 + m^2)}{6},$$

$$c = \frac{(2 - m^2)a}{3}, \beta = \frac{a}{12}, \quad (10)$$

$$v_7 = a \operatorname{sn}(b\xi)^2, A = -\frac{a^2}{6m^2},$$

$$c = \frac{(1 + m^2)a}{3m^2}, \beta = -\frac{a}{12b^2m^2}, \quad (11)$$

$$v_8 = a \operatorname{cn}(b\xi)^2, A = \frac{a^2(1 - m^2)}{6m^2},$$

$$c = \frac{(2m^2 - 1)a}{3m^2}, \beta = \frac{a}{12b^2m^2}, \quad (12)$$

$$v_9 = a \operatorname{dn}(b\xi)^2, A = \frac{a^2(-1 + m^2)}{6},$$

$$c = \frac{(2 - m^2)a}{3}, \beta = \frac{a}{12b^2}, \quad (13)$$

式中 a, b 为实常数, $\operatorname{sn}, \operatorname{cn}, \operatorname{dn}$ 为 Jacobi 椭圆函数, m 为模数.

2) 当 $y_1 > y_2 > y_3, a > 0$, 且均为实常数, (4) 式的 v 有下列 2 组 Jacobi 椭圆函数解:

$$v_{10} = y_3 + (y_2 - y_3) \operatorname{sn}\left(\frac{\sqrt{a(y_1 - y_3)}\xi}{2}\right)^2,$$

$$A = -\frac{y_1 y_2 + y_1 y_3 + y_3 y_2}{6},$$

$$B = \frac{y_1 y_2 y_3}{6},$$

$$c = \frac{y_1 + y_2 + y_3}{3},$$

$$\beta = -\frac{1}{3a},$$

$$m = \sqrt{\frac{y_2 - y_3}{y_1 - y_3}}; \quad (14)$$

$$v_{11} = y_2 + (y_1 - y_2) \operatorname{cn}\left(\frac{\sqrt{a(y_1 - y_3)}\xi}{2}\right)^2,$$

$$A = -\frac{y_1 y_2 + y_1 y_3 + y_3 y_2}{6},$$

$$B = \frac{y_1 y_2 y_3}{6},$$

$$c = \frac{y_1 + y_2 + y_3}{3},$$

$$\beta = \frac{1}{3a},$$

$$m = \sqrt{\frac{y_1 - y_2}{y_1 - y_3}}. \quad (15)$$

3) 当 $A = 0, B = 0, c = \frac{a}{3b}, \beta = -\frac{1}{3b}$ (4) 式的 v 有双曲函数和三角函数解:

$$v_{12} = -\frac{a \operatorname{sech}\left(\frac{\sqrt{a}\xi}{2}\right)^2}{b}, a > 0, b < 0; \quad (16)$$

$$v_{13} = \frac{a \operatorname{csch}\left(\frac{\sqrt{a}\xi}{2}\right)^2}{b}, a > 0, b > 0; \quad (17)$$

$$v_{14} = -\frac{a \operatorname{sec}\left(\frac{\sqrt{-a}\xi}{2}\right)^2}{b}, a < 0. \quad (18)$$

式中 a, b 为实常数.

4) 当 $c = 0$, (4) 式的 v 有 Weierstrass 椭圆函数解:

$$v_{15} = P(\xi, a, b), A = \frac{a}{24}, B = \frac{b}{24}, \beta = -\frac{1}{12}, \quad (19)$$

式中 a, b 为 Weierstrass 椭圆函数不变量.

为了利用 (4) 式作变换, 将 (4) 式的 v 对 ξ 求导, 得

$$\frac{d^2 v}{d\xi^2} = \frac{-v(\xi)^2 + 2c v(\xi) + 2A}{2\beta},$$

$$\frac{d^3 v}{d\xi^3} = \sqrt{3} \sqrt{\frac{-v(\xi)^3 + 3c v(\xi)^2 + 6A v(\xi) + 6B}{\beta}}$$

$$\times (-v(\xi) + c)(3\beta),$$

$$\frac{d^4 v}{d\xi^4} = (5v(\xi)^3 - 15c v(\xi)^2 - 18A v(\xi) + 6c^2 v(\xi) + 6cA - 12B)(6\beta^2). \quad (20)$$

对于非线性偏微分方程

$$\mathcal{G}(u, u_t, u_{x,t}, u_x, u_{x,x}, \dots, F) = 0, \quad (21)$$

其中 F 是依赖于时间 t 的函数, 为外力项. 设 (21) 式有如下形式的解:

$$u = P + Qv(\xi), \xi = \xi(x, t), P = P(x, t),$$

$$Q = Q(x, t), \quad (22)$$

式中 $v(\xi)$ 为 (4) 式的解. 将 (22) 式代入 (21) 式, 得到关于 $\frac{d^i v}{d\xi^i}$ 的方程, 用 (4) 和 (20) 式右边替换左边, 合并

有关 $v(\xi) \left(\sqrt{\frac{-3u(\xi)^3 + 9u(\xi)^2 + 18Au(\xi) + 18B}{\beta}} \right)^k$ 的

同类项, 并令它们的系数为零, 从而可定得 P, Q, ξ . 我们把这一方法, 应用于求解含外力项的变系数广义 KdV 方程.

3. 变系数广义 KdV 方程 (1) 式的解

首先, 将 (22) 式代入 (1) 式, 得到关于 $\frac{d^i v}{d\xi^i} (i = 1, 2, 3, 4)$ 的方程, 用 (4) 和 (20) 式右边替换 $\frac{d^i v}{d\xi^i}$, 合并

$v(\xi) \left(\sqrt{\frac{-3u(\xi)^3 + 9u(\xi)^2 + 18Au(\xi) + 18B}{\beta}} \right)^k$ 的

同类项, 当 $k=0, j=0, 1, \dots, A$; 当 $k=1, j=0, 1$, 并令它的系数为零, 由此可得到如下的偏微分方程组:

$$\begin{aligned} & (6F_1 PP_{x,x}\beta^2 - 2F_2 Q\xi_x^4 B + F_2 Q\xi_x^4 cA \\ & + Q\xi_x \xi_{x,x} \beta A + 6F_1 PQ\xi_x^2 \beta A + 12F_2 Q_x \xi_x \xi_{x,x} \beta A \\ & + 12F_1 Q^2 \xi_x^2 \beta B + 3F_2 Q\xi_x^2 \beta A + 6F_2 Q_x \xi_x^2 \beta A \\ & + 4F_2 Q\xi_x \xi_{x,x} \beta A + P_{x,x} \beta^2 + F_4 \beta^2 \\ & + F_2 P_{x,x,x} \beta^2 + 6F_1 P_{x,x}^2 \beta^2 \\ & + 6F_3 P_x \beta^2) \beta^2 = 0, \end{aligned} \quad (23)$$

$$\begin{aligned} & (F_2 Q_{x,x,x} \beta^2 + 6F_3 Q_x \beta^2 + 6F_1 PQ_x \beta^2 \\ & + 12F_1 P_x Q_x \beta^2 + 6F_1 QP_{x,x} \beta^2 - 3F_2 Q\xi_x^4 A \\ & + F_2 Q\xi_x^4 c^2 + Q\xi_x \xi_{x,x} \beta c + 18F_1 Q^2 \xi_x^2 \beta A + 6F_1 PQ\xi_x^2 \beta c \\ & + 12F_2 Q_x \xi_x \xi_{x,x} \beta c + 3F_2 Q\xi_x^2 \beta c + 6F_2 Q_x \xi_x^2 \beta c \\ & + 4F_2 Q\xi_x \xi_{x,x} \beta c + Q_{x,x} \beta^2) \beta^2 = 0, \end{aligned} \quad (24)$$

$$\begin{aligned} & (12F_1 Q_x^2 \beta^2 + 12F_1 QQ_{x,x} \beta^2 - 5F_2 Q\xi_x^4 c \\ & - Q\xi_x \xi_{x,x} \beta + 24F_1 Q^2 \xi_x^2 \beta c - 6F_1 PQ\xi_x^2 \beta \\ & - 12F_2 Q_x \xi_x \xi_{x,x} \beta - 3F_2 Q\xi_x^2 \beta \\ & - 6F_2 Q_x \xi_x^2 \beta - 4F_2 Q\xi_x \xi_{x,x} \beta) (2\beta^2) = 0, \end{aligned} \quad (25)$$

$$- \frac{5\xi_x^2 Q(-F_2 \xi_x^2 + 6F_1 Q\beta)}{6\beta^2} = 0, \quad (26)$$

$$\begin{aligned} & (Q_x \xi_x \beta + Q_x \xi_x \beta + Q_{x,x} \beta + F_2 Q\xi_{x,x} \beta \\ & + 6F_3 Q\xi_x \beta + 4F_2 Q_x \xi_{x,x} \beta + 6F_2 Q_x \xi_x \beta \\ & + 4F_2 Q_{x,x} \xi_x \beta + 12F_1 P_x Q\xi_x \beta + 12F_1 PQ_x \xi_x \beta \\ & + 6F_1 PQ\xi_{x,x} \beta + 4F_2 Q_x \xi_x^3 c \\ & + 6F_2 Q\xi_x^2 \xi_{x,x} c) (3\beta) = 0, \end{aligned} \quad (27)$$

$$\begin{aligned} & (3F_1 Q^2 \xi_{x,x} \beta + 12F_1 Q_x Q\xi_x \beta - 2F_2 Q_x \xi_x^3 \\ & - 3F_2 Q\xi_x^2 \xi_{x,x} c) (3\beta) = 0, \end{aligned} \quad (28)$$

式中 $P_{x,t} = \frac{\partial^2 P}{\partial x \partial t}, Q_{x,t} = \frac{\partial^2 Q}{\partial x \partial t}, \xi_{x,t} = \frac{\partial^2 \xi}{\partial x \partial t}, \dots$ 由 (26) 式解得

$$Q = \frac{F_2 \xi_x^2}{6F_1 \beta}, \quad (29)$$

将 (29) 式分别代入 (23) (24) (25) (27) 和 (28) 式, 解得

$$\xi = F_5 x + F_6, P = - \frac{F_2 F_5^3 c + \beta F_{5,t} x + \beta F_{6,t}}{6F_5 F_1 \beta},$$

$$Q = \frac{F_2 F_5^2}{6F_1 \beta}, \quad (30)$$

式中 F_6 为 t 的任意函数, F_1, F_2, F_3, F_4 和 F_5 是受下列条件制约的 t 的函数:

$$F_3 = - \frac{F_{2,t} F_5 F_1 + F_2 F_{5,t} F_1 - F_2 F_5 F_{1,t}}{6F_2 F_5 F_1} \quad (31)$$

$$F_4 = - \frac{3F_{5,t}^2 F_2 - F_2 F_{5,t,t} F_5 + F_{5,t} F_{2,t} F_5}{6F_1 F_5^2 F_2} \quad (32)$$

将 (30) 式代入 (22) 式, 定得满足 (1) 式新的解为

$$u = - \frac{F_2 F_5^3 c + \beta F_{5,t} x + \beta F_{6,t}}{6F_5 F_1 \beta} + \frac{F_2 F_5^3 v}{6F_1 \beta}, \quad (33)$$

式中 $F_{1,t} = \frac{dF_1}{dt}, F_{2,t} = \frac{dF_2}{dt}, \dots, F_{6,t} = \frac{dF_6}{dt}$, 用 (5) —

(19) 式中的 $v_1, v_2, \dots, v_{15}, c, \beta$ 分别代替 (33) 式中的 c, v, β , 可得到 15 组满足 (1) 式新的精确解:

$$u_1 = - \frac{-4F_2 F_5^3 - 4F_2 F_5^3 m^2 + F_{5,t} x + F_{6,t} + 12F_2 F_5^3 m^2 \text{sr}(\xi)}{6F_5 F_1},$$

$$u_2 = \frac{-8F_2 F_5^3 m^2 + 4F_2 F_5^3 - F_{5,t} x - F_{6,t} + 12F_2 F_5^3 m^2 \text{cr}(\xi)}{6F_5 F_1},$$

$$u_3 = \frac{-8F_2 F_5^3 + 4F_2 F_5^3 m^2 - F_{5,t} x - F_{6,t} + 12F_2 F_5^3 \text{dr}(\xi)}{6F_5 F_1},$$

$$u_4 = \frac{-8F_2 F_5^3 m^2 + 4F_2 F_5^3 - F_{5,t} x - F_{6,t} + 12F_2 F_5^3 m^2 \text{sr}(\xi)}{6F_5 F_1},$$

$$u_5 = \frac{-8F_2 F_5^3 m^2 + 4F_2 F_5^3 - F_{5,t} x - F_{6,t} + 12F_2 F_5^3 m^2 \text{cr}(\xi)}{6F_5 F_1},$$

$$u_6 = \frac{-8F_2F_5^3 + 4F_2F_5^3m^2 - F_{5,\mu}x - F_{6,\mu} + 12F_2F_5^3\operatorname{dn}(\xi)^2}{6F_5F_1},$$

$$u_7 = -(-4F_2F_5^3b^2 - 4F_2F_5^3b^2m^2 + F_{5,\mu}x + F_{6,\mu} + 12F_2F_5^3b^2m^2\operatorname{sn}(b\xi)^2)(6F_5F_1),$$

$$u_8 = (-8F_2F_5^3b^2m^2 + 4F_2F_5^3b^2 - F_{5,\mu}x - F_{6,\mu} + 12F_2F_5^3b^2m^2\operatorname{cn}(b\xi)^2)(6F_5F_1),$$

$$u_9 = (-8F_2F_5^3b^2 + 4F_2F_5^3b^2m^2 - F_{5,\mu}x - F_{6,\mu} + 12F_2F_5^3b^2\operatorname{dn}(b\xi)^2)(6F_5F_1),$$

$$u_{10} = -\left(-F_2F_5^3ay_1 - F_2F_5^3ay_2 + 2F_2F_5^3ay_3 + F_{5,\mu}x + F_{6,\mu} - 3F_2F_5^3\alpha(y_2 - y_3)\operatorname{sn}\left(\frac{\sqrt{\alpha(y_1 - y_3)}\xi}{2}\right)^2\right)(6F_5F_1),$$

$$u_{11} = \left(-F_2F_5^3ay_1 + 2F_2F_5^3ay_2 - F_2F_5^3ay_3 - F_{5,\mu}x - F_{6,\mu} - 3F_2F_5^3\alpha(y_1 - y_2)\operatorname{cn}\left(\frac{\sqrt{\alpha(y_1 - y_3)}\xi}{2}\right)^2\right)(6F_5F_1),$$

$$u_{12} = \frac{-F_2F_5^3a - F_{5,\mu}x - F_{6,\mu} + 3F_2F_5^3a\operatorname{sech}\left(\frac{\sqrt{a}\xi}{2}\right)^2}{6F_5F_1},$$

$$u_{13} = \frac{-F_2F_5^3a - F_{5,\mu}x - F_{6,\mu} - 3F_2F_5^3a\operatorname{csch}\left(\frac{\sqrt{a}\xi}{2}\right)^2}{6F_5F_1},$$

$$u_{14} = \frac{-F_2F_5^3a - xF_{5,\mu} - F_{6,\mu} + 3F_2F_5^3a\operatorname{sec}\left(\frac{\sqrt{-a}\xi}{2}\right)^2}{6F_5F_1},$$

$$u_{15} = -\frac{F_{5,\mu}x + F_{6,\mu} + 12F_2F_5^3\rho(\xi, a, b)}{6F_5F_1}.$$

式中 ξ 由 (30) 式决定. 其中 u_{12}, u_{13} 为类孤波解, u_{15} 为 Weierstrass 椭圆函数解. 当 $m \rightarrow 1$, $\operatorname{sn}(\xi) = \tanh(\xi)$, $\operatorname{cn}(\xi) = \operatorname{sech}(\xi)$, $\operatorname{dn}(\xi) = \operatorname{sech}(\xi)$, $u_1 \dots u_{11}$ 由 Jacobi 椭圆函数精确解退化为类孤波解; 当 $m \rightarrow 0$, $\operatorname{sn}(\xi) = \sin(\xi)$, $\operatorname{dn}(\xi) = 1$, $\operatorname{cn}(\xi) = \cos\xi$, $u_1 \dots u_{11}$ 由 Jacobi 椭圆函数精确解退化为三角函数精确解, 不再列出. 由于解 (33) 式不含 A, B , 所以 $u_2 = u_5$, $u_3 = u_6$.

4. 方程 (2) 式的解

如果取 (30) 式中的 $F_5 = k_1$, k_1 为常数, 由 (30)–(32) 式得

$$P = -\frac{F_2k_1^3c + \beta F_{6,\mu}}{6k_1F_1\beta}, Q = \frac{F_2k_1^2}{6F_1\beta}, F_4 = 0,$$

$$F_3 = -\frac{F_{2,\mu}F_1 - F_2F_{1,\mu}}{6F_2F_1}, \xi = k_1x + F_6. \quad (34)$$

将 (34) 式代入 (22) 式, 定得满足 (2) 式新的精确解为

$$u = -\frac{F_2k_1^3c + \beta F_{6,\mu} - F_2k_1^3u(\xi)}{6k_1F_1\beta}. \quad (35)$$

用 (5)–(19) 式中的 $v_1, v_2, \dots, v_{15}, c, \beta$ 分别代替 (35) 式中的 $c, \rho(\xi), \beta$, 得到 15 组满足 (2) 式的新的精确解、类孤波解、三角函数解和 Weierstrass 椭圆函数解. 式中 ξ 由 (34) 式决定. 解的情况类似于 3.

5. 结 论

方程 (1) 比 (3) 复杂, 我们相信一个复杂系统的物理行为包含着简单系统的物理行为, 简单系统的行为是复杂系统行为的特例, 简单系统的解与复杂系统的解, 存在一定的联系. 基于这一想法, 本文用普通 KdV 方程的 (4) 式作变换, 构造含外力项变系数广义 KdV 方程的解, 获得了 (1) 和 (2) 式新的精确解、类孤波解、三角函数解和 Weierstrass 椭圆函数解. 解的结果, 用 Maple 进行了验算. 而本方法是否具有普遍的意义, 能否应用于构造含外力项变系数广义 (2+1) 维和 (3+1) 维的 KdV 方程的解, 我们将作进一步的研究.

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New solitary-wave-like solutions and exact solutions to variable coefficient generalized KdV equation

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Abstract

By the transforming with general KdV equation , the solutions of the variable coefficient generalized KdV equation are constructed firstly. As a result , we successfully obtain the new exact Jacobi elliptic function solutions , solitary-wave-like solutions , trigonometric function solutions and Weierstrass elliptic function solutions of variable coefficient generalized KdV equation.

Keywords : KdV equation , variable coefficient generalized KdV equation , solitary-wave-like solution , exact solution

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