

# 直立防波堤上部分反射的三阶双色双向水波理论\*

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为充分克服以往直立防波堤上立波波浪力的内在局限性和刻画直立防波堤堤前的多色多向波, 提出一种可包含经典二维长峰波解和三维单色单向短峰波解的部分反射的三阶双色双向水波理论, 并从中推断出适用于高阶多色多向波系统的一项普适法则——偶数阶频率为零, 奇数阶频率个数构成一等比数列.

**关键词:** 双色双向水波, 三阶理论, 普适法则, 部分反射波

**PACC:** 0340K, 0200, 9210H

## 1. 引 言

波浪力为近海、海洋建筑物上所遭受的一种主要荷载, 在当今全球急剧发展的海洋工程实践中至关重要. 如何确定最为常见的直立防波堤上的波浪力呢? 如果堤前处于非破碎波状态, 通常就以理想、简单的波浪正入射直立堤上而发生全反射, 最终导致的二维长峰立波波浪力为标准. 这里, 存在着方便、实用甚至可靠的特色, 但其内在缺陷是明显的. 例如, 立波波浪力确实为“最大、最危险”的荷载吗? 这并非想象或简单的理论、实验所能确定.

现在, 面对着直立堤, 可以明确以下两个基本物理事实: 堤前波浪一般并非处于假定的正入射和全反射状态, 而是通常的倾斜入射、倾斜部分反射; 入射波、反射波并非受制于理想的单色单向限制, 理应是多色多向的. 如果依据以上事实而得以谋划一理论, 则其包容性、先进性是明确和必要的.

早在三十年前, Hsu 等<sup>[1]</sup>就提出了一种可发生于直立堤前的经典三阶三维单色单向全反射短峰波理论. 正是借助该理论, Fenton<sup>[2]</sup>后来发现: 倾斜入射直立堤时而产生的波浪力显著大于相应的立波波浪力. 至今, 短峰波已在多个方面得到大大拓展<sup>[3-8]</sup>. 其中就包括以两列自由交叉波传播为背景而于最近构造的三阶双色双向波理论<sup>[6]</sup>. 显见, 目

前这种多色多向波理论发展远远落后于二维长峰波<sup>[8-10]</sup>和三维短峰波的进程. 其原因就在于前者理论构造的极大繁杂性. 若要真正反映一般或典型工程物理场景多色多向波运动的本质属性, 则从双色双向波入手, 就不失为一个实际的良好开端.

基于以上考虑, 本文提出一种可比较实际地考虑直立堤前非破碎波入射和部分反射波态的三阶双色双向波理论, 以此考察、确定多色多向波的一般结构特征.

## 2. 构 造

假定流体做无黏、不可压缩的无旋流体运动, 其速度势为  $\phi(x, y, z, t)$ . 采用如图 1 所示的直角坐标系:  $x$  轴沿着直立防波堤方向,  $y$  轴垂直于堤,  $z$  轴竖直向上, 原点位于静水面上.  $z = \zeta(x, y, t)$  和  $z = -h$  分别代表自由表面和常水深海底,  $y = 0$  则表示直立堤. 假定两列波分别以与  $y$  轴成  $\theta_m, \theta_n$  的角度入射、反射, 其角频率分别为  $\omega_m, \omega_n$ , 并且入射波或反射波的平均波数和平均频率分别为  $\bar{k}$  和  $\bar{\omega}$ , 以此可引入如下无量纲化量:

$$\begin{aligned} & (\bar{x}, \bar{y}, \bar{z}, \bar{\zeta}, \bar{h}) \\ & = \bar{k}(x, y, z, \zeta, h), \quad \bar{\phi} = \phi \sqrt{\frac{\bar{k}^3}{g}}, \end{aligned}$$

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$$\begin{aligned} \bar{t} &= \omega t, \bar{\omega} = \frac{\omega}{\sqrt{gk}} = \frac{\omega}{\omega}, \\ \bar{k}_i &= \frac{k_i}{k}, \quad \bar{\omega}_i = \frac{\omega_i}{\omega}, \end{aligned} \quad (1)$$

其中,  $\omega$  和  $g$  分别代表角频率和重力加速度,  $k_i = \frac{2\pi}{L_i}$  ( $i = m, n$ ), 并且  $k_i$  的分量可表示为

$$\begin{aligned} k_{ix} &= \frac{2\pi}{L_{ix}} = k_i \sin \theta_i = k_i i_1, \\ k_{iy} &= \frac{2\pi}{L_{iy}} = k_i \cos \theta_i = k_i i_2. \end{aligned} \quad (2)$$

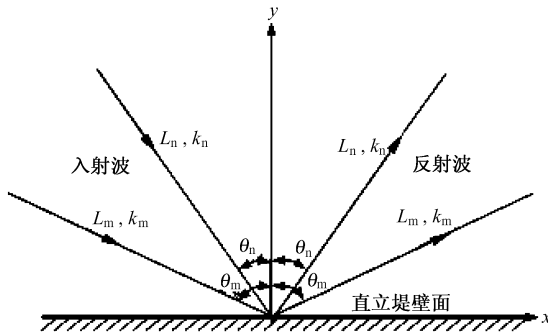


图1 双色双向波入射、反射直立堤

现在, 可将短峰波的控制方程及其边界条件无量纲化为(省略无量纲符号)

$$\nabla^2 \phi = 0 \quad (-h \leq z \leq \zeta), \quad (3)$$

$$\zeta + \omega \phi_t + \frac{1}{2}(\phi_x^2 + \phi_y^2 + \phi_z^2) = 0 \quad (z = \zeta), \quad (4)$$

$$\phi_z - \omega \zeta_t - (\phi_x \zeta_x + \phi_y \zeta_y) = 0 \quad (z = \zeta), \quad (5)$$

$$\phi_z = 0 \quad (z = -h), \quad (6)$$

其中,  $\nabla = \left( \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$ ,  $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$ .

将  $\phi, \zeta$  和  $\omega$  进行如下摄动展开:

$$\begin{aligned} \phi &= \varepsilon \phi_1 + \varepsilon^2 \phi_2 + \dots, \\ \zeta &= \varepsilon \zeta_1 + \varepsilon^2 \zeta_2 + \dots, \\ \omega &= \omega_0 + \varepsilon \omega_1 + \varepsilon^2 \omega_2 + \dots. \end{aligned} \quad (7)$$

将自由表面速度势在  $z = 0$  处进行 Taylor 展开, 可得下列各阶 ( $r = 1, 2, 3$ ) 控制方程组, 从而可依次求解.

$$\nabla^2 \phi_r = 0 \quad (-h \leq z \leq 0), \quad (8)$$

$$\zeta_r + \omega_0 \phi_{rt} = E_r \quad (z = 0), \quad (9)$$

$$\phi_{rz} + \omega_0^2 \phi_{rtt} = F_r \quad (z = 0), \quad (10)$$

$$\phi_{rz} = 0 \quad (z = -h), \quad (11)$$

其中,

$$E_1 = F_1 = 0, \quad (12)$$

$$E_2 = -\omega_1 \phi_{1t} - \omega_0 \phi_{1tz} \zeta_1 - \frac{1}{2}(\phi_{1x}^2 + \phi_{1y}^2 + \phi_{1z}^2), \quad (13)$$

$$F_2 = -2\omega_0 \omega_1 \phi_{1tt} - 2\omega_0 \nabla \phi_1 \cdot \nabla \phi_{1t} - \zeta_1 (\phi_{1zz} + \omega_0^2 \phi_{1ztt}), \quad (14)$$

$$\begin{aligned} E_3 &= -\omega_2 \phi_{1t} - \omega_1 (\phi_{2t} + \phi_{1zt} \zeta_1) \\ &\quad - \omega_0 (\phi_{1zt} \zeta_2 + \phi_{2zt} \zeta_1 + \frac{1}{2} \phi_{1ztt} \zeta_1^2) \\ &\quad - \nabla \phi_1 \cdot \nabla \phi_2 - \zeta_1 \nabla \phi_1 \cdot \nabla \phi_{1z}, \end{aligned} \quad (15)$$

$$\begin{aligned} F_3 &= -2\omega_2 \omega_0 \phi_{1tt} \\ &\quad - 2\omega_1 [\omega_0 (\phi_{2tt} + \phi_{1ztt} \zeta_1) + \nabla \phi_1 \cdot \nabla \phi_{1t}] \\ &\quad - 2\omega_0 [\nabla \phi_1 \cdot \nabla \phi_{2t} + \nabla \phi_2 \cdot \nabla \phi_{1t} \\ &\quad + \zeta_1 (\nabla \phi_{1z} \cdot \nabla \phi_{1t} + \nabla \phi_1 \cdot \nabla \phi_{1zt})] \\ &\quad - \omega_1^2 \phi_{1tt} - \omega_0^2 (\zeta_1 \phi_{2ztt} + \zeta_2 \phi_{1ztt} + \frac{1}{2} \zeta_1^2 \phi_{1ztt}) \\ &\quad - (\zeta_1 \phi_{2zz} + \zeta_2 \phi_{1zz} + \frac{1}{2} \zeta_1^2 \phi_{1zz}) \\ &\quad - \frac{1}{2} \nabla \phi_1 \cdot \nabla [(\nabla \phi_1)^2]. \end{aligned} \quad (16)$$

### 3. 三阶解

#### 3.1. 第一阶解

$$\begin{aligned} \zeta_1 &= a_m \cos \alpha_m + b_m \cos \beta_m + a_n \cos \alpha_n \\ &\quad + b_n \cos \beta_n, \end{aligned} \quad (17)$$

$$\begin{aligned} \phi_1 &= \frac{\cosh k_m Z}{\omega_m \cosh k_m h} (a_m \sin \alpha_m + b_m \sin \beta_m) \\ &\quad + \frac{\cosh k_n Z}{\omega_n \cosh k_n h} (a_n \sin \alpha_n + b_n \sin \beta_n) \\ &\quad (Z = z + h), \end{aligned} \quad (18)$$

$$\omega_0^2 = \tanh h, \quad (19)$$

其中,

$$\alpha_i = k_i (i_1 x - i_2 y) - \frac{\omega_i t}{\omega_0},$$

$$\beta_i = k_i (i_1 x + i_2 y) - \frac{\omega_i t}{\omega_0},$$

$$\omega_i^2 = k_i \tanh k_i h. \quad (20)$$

显见, 第一阶解呈一对称形式, 两列波各自独立地线性叠加, 蕴含着单色单向短峰波第一阶解<sup>[1]</sup>. 为确定该短峰波高阶解的某些系数而针对特定二维立波运动所得到的若干附加条件<sup>[11]</sup>, 应该也

是确定该双色双向波高阶解相应某些系数的附加条件. 这里只需对这些附加条件中的振幅方程做出局部修正, 即

$$\int_{-h}^0 \int_0^\pi \int_0^{2\pi} \phi_s(y, z, t) \sin t \cos y \, dt \, dy \, dz = -2c_s \pi^2 (\tanh h)^{1/2}, \quad (21)$$

其中,  $\phi_s$  和  $c_s$  依次表征立波的速度势和振幅. 最后, 为保证该双色双向波高阶解的唯一性, 需满足如下条件<sup>[11]</sup>:

$$\frac{p \tanh p h}{\tanh h} \neq j^2 \quad (p \geq 2, j \geq 1). \quad (22)$$

### 3.2. 第二阶解

$$\begin{aligned} \zeta_2 = & G_{2m} (a_m^2 \cos 2\alpha_m + b_m^2 \cos 2\beta_m) \\ & + G_{2n} (a_n^2 \cos 2\alpha_n + b_n^2 \cos 2\beta_n) \\ & + G_{mm}^+ a_m b_m \cos(\alpha_m + \beta_m) \\ & + G_{nn}^+ a_n b_n \cos(\alpha_n + \beta_n) \\ & + G_{mm}^- a_m b_m \cos(\alpha_m - \beta_m) \\ & + G_{nn}^- a_n b_n \cos(\alpha_n - \beta_n) \\ & + G_{mn}^+ [a_m a_n \cos(\alpha_m + \alpha_n) \\ & + b_m b_n \cos(\beta_m + \beta_n)] \\ & + G_{mn}^- [a_m a_n \cos(\alpha_m - \alpha_n) \\ & + b_m b_n \cos(\beta_m - \beta_n)] \\ & + G_{nm}^+ [a_n b_m \cos(\alpha_n + \beta_m) \\ & + b_n a_m \cos(\beta_n + \alpha_m)] \\ & + G_{nm}^- [a_n b_m \cos(\alpha_n - \beta_m) \\ & + b_n a_m \cos(\beta_n - \alpha_m)], \end{aligned} \quad (23)$$

$$\begin{aligned} \phi_2 = & \beta_0 + \alpha_0 t + F_{2m} (a_m^2 \sin 2\alpha_m \\ & + b_m^2 \sin 2\beta_m) \cosh k_{2m} Z \\ & + F_{2n} (a_n^2 \sin 2\alpha_n + b_n^2 \sin 2\beta_n) \cosh k_{2n} Z \\ & + F_{mm}^+ a_m b_m \sin(\alpha_m + \beta_m) \cosh k_{mm} Z \\ & + F_{nn}^+ a_n b_n \sin(\alpha_n + \beta_n) \cosh k_{nn} Z \\ & + F_{mn}^+ [a_m a_n \sin(\alpha_m + \alpha_n) \\ & + b_m b_n \sin(\beta_m + \beta_n)] \cosh k_{mn}^+ Z \\ & + F_{mn}^- [a_m a_n \sin(\alpha_m - \alpha_n) \\ & + b_m b_n \sin(\beta_m - \beta_n)] \cosh k_{mn}^- Z \\ & + F_{nm}^+ [a_n b_m \sin(\alpha_n + \beta_m) \\ & + b_n a_m \sin(\beta_n + \alpha_m)] \cosh k_{nm}^+ Z \\ & + F_{nm}^- [a_n b_m \sin(\alpha_n - \beta_m) \\ & + b_n a_m \sin(\beta_n - \alpha_m)] \cosh k_{nm}^- Z, \end{aligned} \quad (24)$$

$$\omega_1 = 0, \quad (25)$$

其中,  $\beta_0$  为一常量. 各传递函数和有关的波数如下 (对于具有同等表达形式的两个对称项, 仅列出涉及下标  $m$  的项):

$$\alpha_0 = \frac{G_{00}^m (a_m^2 + b_m^2) + G_{00}^n (a_n^2 + b_n^2)}{\omega_0}, \quad (26)$$

$$G_{00}^m = A_{21}(\omega_m, \gamma_m^+), \quad (27)$$

$$G_{2m} = A_{22}(\omega_m, \gamma_m^+) + 2\omega_m F_{2m} \cosh k_{2m} h, \quad (28)$$

$$G_{mm}^+ = A_{22}(\omega_m, \gamma_m^-) + 2\omega_m F_{2m} \cosh k_{2m} h, \quad (29)$$

$$G_{mm}^- = A_{21}(\omega_m, \gamma_m^-), \quad (30)$$

$$\begin{aligned} G_{mn}^\pm = & A_{23}(\omega_m, \pm \omega_n, \gamma_{mn}^\pm) \\ & + (\omega_m \pm \omega_n) F_{mn}^\pm \cosh k_{mn}^\pm h, \end{aligned} \quad (31)$$

$$\begin{aligned} G_{nm}^\pm = & A_{23}(\pm \omega_n, \omega_m, \gamma_{nm}^\pm) \\ & + (\omega_n \pm \omega_m) F_{nm}^\pm \cosh k_{nm}^\pm h, \end{aligned} \quad (32)$$

$$F_{2m} = \frac{A_{2m}}{2k_{2m} \sinh k_{2m} h - 8\omega_m^2 \cosh k_{2m} h}, \quad (33)$$

$$F_{mm}^+ = \frac{A_{mm}^+}{k_{mm}^+ \sinh k_{mm}^+ h - 4\omega_m^2 \cosh k_{mm}^+ h}, \quad (34)$$

$$F_{mn}^\pm = \frac{A_{mn}^\pm}{k_{mn}^\pm \sinh k_{mn}^\pm h - (\omega_m \pm \omega_n)^2 \cosh k_{mn}^\pm h}, \quad (35)$$

$$F_{nm}^\pm = \frac{A_{nm}^\pm}{k_{nm}^\pm \sinh k_{nm}^\pm h - (\omega_n \pm \omega_m)^2 \cosh k_{nm}^\pm h}, \quad (36)$$

$$k_{2m} = 2k_m, \quad (37)$$

$$k_{mm}^+ = 2k_m m_1, \quad (38)$$

$$k_{mn}^\pm = \sqrt{k_m^2 + k_n^2 \pm 2\gamma_{mn}^\pm}, \quad (39)$$

$$k_{nm}^\pm = \sqrt{k_m^2 + k_n^2 \pm 2\gamma_{nm}^\pm}, \quad (40)$$

其中, 所引入的函数为

$$A_{2m} = \Gamma_{22}(\omega_m, k_m, \gamma_m^+), \quad (41)$$

$$A_{mm}^+ = \Gamma_{22}(\omega_m, k_m, \gamma_m^-), \quad (42)$$

$$\Gamma_{22}(\omega_m, k_m, \gamma_m^+) = \frac{3\omega_m^4 - k_m^2 - 2\gamma_m^+}{\omega_m}, \quad (43)$$

$$\gamma_m^\pm = k_m^2 (m_1^2 \pm m_2^2), \quad (44)$$

$$A_{mn}^\pm = \Gamma_{23}(\omega_m, \pm \omega_n, k_m, \pm k_n, \pm \gamma_{mn}^\pm), \quad (45)$$

$$A_{nm}^\pm = \pm \Gamma_{23}(\pm \omega_n, \omega_m, \pm k_n, k_m, \pm \gamma_{nm}^\pm), \quad (46)$$

$$\begin{aligned} & \Gamma_{23}(\omega_m, \omega_n, k_m, k_n, \gamma_{mn}^+) \\ & = \omega_m \omega_n (\omega_m + \omega_n) + \frac{1}{2} \left( \frac{\omega_m^4 - k_m^2}{\omega_m} + \frac{\omega_n^4 - k_n^2}{\omega_n} \right) \\ & \quad - \gamma_{mn}^+ \left( \frac{1}{\omega_m} + \frac{1}{\omega_n} \right), \end{aligned} \quad (47)$$

$$\gamma_{mn}^* = \gamma_{nm}^* = k_m k_n (m_1 n_1 \pm m_2 n_2), \quad (48)$$

$$A_{21}(\omega_m, \gamma_m^+) = \frac{\omega_m^4 - \gamma_m^+}{4\omega_m^2}, \quad (49)$$

$$A_{22}(\omega_m, \gamma_m^+) = \frac{3\omega_m^4 - \gamma_m^+}{4\omega_m^2}, \quad (50)$$

$$A_{23}(\omega_m, \omega_n, \gamma_{mn}^+) = \frac{1}{2} \left( \omega_m^2 + \omega_n^2 + \omega_m \omega_n - \frac{\gamma_{mn}^+}{\omega_m \omega_n} \right). \quad (51)$$

可以看出, 第二阶解—— $\zeta_2$  与  $\phi_2$  相互对称地存在对应项. 通过四种基本位相  $\alpha_m, \beta_m, \alpha_n, \beta_n$ , 两列波以下述四种形式组合在一起: 同波同相和:  $2\alpha_m, 2\beta_m, 2\alpha_n, 2\beta_n$ ; 同波异相和差:  $\alpha_m \pm \beta_m, \alpha_n \pm \beta_n$ ; 异波同函数相和差:  $\alpha_m \pm \alpha_n, \beta_m \pm \beta_n$ ; 异波异函数相和差:  $\alpha_n \pm \beta_m; \alpha_m \pm \beta_n$ .

实际上, 这四种基本的第二阶位相组合可预先判知, 并可称为对称性、位相决定第二阶波.

### 3.3. 第三阶解

$$\begin{aligned} \zeta_3 = & G_{m0} a_m \cos \alpha_m + G_{n0} a_n \cos \alpha_n + G_{0m} b_m \cos \beta_m \\ & + G_{0n} b_n \cos \beta_n + G_{3m} (a_m^3 \cos 3\alpha_m + b_m^3 \cos 3\beta_m) \\ & + G_{3n} (a_n^3 \cos 3\alpha_n + b_n^3 \cos 3\beta_n) \\ & + G_{m2m}^- [a_m b_m^2 \cos(\alpha_m - 2\beta_m) \\ & + b_m a_m^2 \cos(\beta_m - 2\alpha_m)] \\ & + G_{n2n}^- [a_n b_n^2 \cos(\alpha_n - 2\beta_n) \\ & + b_n a_n^2 \cos(\beta_n - 2\alpha_n)] \\ & + G_{2mm}^+ [a_m^2 b_m \cos(2\alpha_m + \beta_m) \\ & + b_m^2 a_m \cos(2\beta_m + \alpha_m)] \\ & + G_{2nn}^+ [a_n^2 b_n \cos(2\alpha_n + \beta_n) \\ & + b_n^2 a_n \cos(2\beta_n + \alpha_n)] \\ & + G_{m2n}^+ [a_m a_n^2 \cos(\alpha_m + 2\alpha_n) \\ & + b_m b_n^2 \cos(\beta_m + 2\beta_n)] \\ & + G_{m2n}^- [a_m a_n^2 \cos(\alpha_m - 2\alpha_n) \\ & + b_m b_n^2 \cos(\beta_m - 2\beta_n)] \\ & + G_{2mn}^+ [a_m^2 a_n \cos(2\alpha_m + \alpha_n) \\ & + b_m^2 b_n \cos(2\beta_m + \beta_n)] \\ & + G_{2mn}^- [a_m^2 a_n \cos(2\alpha_m - \alpha_n) \\ & + b_m^2 b_n \cos(2\beta_m - \beta_n)] \\ & + G_{n2m}^+ [a_n b_m^2 \cos(\alpha_n + 2\beta_m) \\ & + b_n a_m^2 \cos(\beta_n + 2\alpha_m)] \end{aligned}$$

$$\begin{aligned} & + G_{n2m}^- [a_n b_m^2 \cos(\alpha_n - 2\beta_m) \\ & + b_n a_m^2 \cos(\beta_n - 2\alpha_m)] \\ & + G_{2nm}^+ [a_n^2 b_m \cos(2\alpha_n + \beta_m) \\ & + b_n^2 a_m \cos(2\beta_n + \alpha_m)] \\ & + G_{2nm}^- [a_n^2 b_m \cos(2\alpha_n - \beta_m) \\ & + b_n^2 a_m \cos(2\beta_n - \alpha_m)] \\ & + G_{mnm}^+ [a_m a_n b_m \cos(\alpha_m + \alpha_n - \beta_m) \\ & + b_m b_n a_m \cos(\beta_m + \beta_n - \alpha_m)] \\ & + G_{mnm}^- [a_m a_n b_m \cos(\alpha_m - \alpha_n - \beta_m) \\ & + b_m b_n a_m \cos(\beta_m - \beta_n - \alpha_m)] \\ & + G_{mnn}^+ [a_m a_n b_n \cos(\alpha_m + \alpha_n - \beta_n) \\ & + b_m b_n a_n \cos(\beta_m + \beta_n - \alpha_n)] \\ & + G_{mnn}^- [a_m a_n b_n \cos(\alpha_m - \alpha_n + \beta_n) \\ & + b_m b_n a_n \cos(\beta_m - \beta_n + \alpha_n)] \\ & + G_{mnn}^+ [a_m b_m b_n \cos(\alpha_m + \beta_m + \beta_n) \\ & + a_m b_m a_n \cos(\alpha_m + \beta_m + \alpha_n)] \\ & + G_{mnn}^- [a_m b_m b_n \cos(\alpha_m + \beta_m - \beta_n) \\ & + a_m b_m a_n \cos(\alpha_m + \beta_m - \alpha_n)] \\ & + G_{nmm}^+ [a_n b_m b_n \cos(\alpha_n + \beta_m) \\ & + \beta_n) + a_n a_m b_n \cos(\alpha_n + \alpha_m + \beta_n)] \\ & + G_{nmm}^- [a_n b_m b_n \cos(\alpha_n - \beta_m + \beta_n) \\ & + a_n a_m b_n \cos(\alpha_n - \alpha_m + \beta_n)], \quad (52) \end{aligned}$$

$$\begin{aligned} \phi_3 = & F_{3m} (a_m^3 \sin 3\alpha_m + b_m^3 \sin 3\beta_m) \cosh 3k_m Z \\ & + F_{3n} (a_n^3 \sin 3\alpha_n + b_n^3 \sin 3\beta_n) \cosh 3k_n Z \\ & + F_{m2m}^- [a_m b_m^2 \sin(\alpha_m - 2\beta_m) \\ & + b_m a_m^2 \sin(\beta_m - 2\alpha_m)] \cosh k_{m2m}^- Z \\ & + F_{n2n}^- [a_n b_n^2 \sin(\alpha_n - 2\beta_n) \\ & + b_n a_n^2 \sin(\beta_n - 2\alpha_n)] \cosh k_{n2n}^- Z \\ & + F_{2mm}^+ [a_m^2 b_m \sin(2\alpha_m + \beta_m) \\ & + b_m^2 a_m \sin(2\beta_m + \alpha_m)] \cosh k_{2mm}^+ Z \\ & + F_{2nn}^+ [a_n^2 b_n \sin(2\alpha_n + \beta_n) \\ & + b_n^2 a_n \sin(2\beta_n + \alpha_n)] \cosh k_{2nn}^+ Z \\ & + F_{m2n}^+ [a_m a_n^2 \sin(\alpha_m + 2\alpha_n) \\ & + b_m b_n^2 \sin(\beta_m + 2\beta_n)] \cosh k_{m2n}^+ Z \\ & + F_{m2n}^- [a_m a_n^2 \sin(\alpha_m - 2\alpha_n) \\ & + b_m b_n^2 \sin(\beta_m - 2\beta_n)] \cosh k_{m2n}^- Z \\ & + F_{2mn}^+ [a_m a_n^2 \sin(2\alpha_m + \alpha_n) \\ & + b_m b_n^2 \sin(2\beta_m + \alpha_n)] \cosh k_{2mn}^+ Z \\ & + F_{2mn}^- [a_m a_n^2 \sin(2\alpha_m - \alpha_n) \\ & + b_m b_n^2 \sin(2\beta_m - \alpha_n)] \cosh k_{2mn}^- Z \end{aligned}$$

$$\begin{aligned}
& + b_m^2 b_n \sin(2\beta_m + \beta_n) ] \cosh k_{2mn}^+ Z \\
& + F_{2mn}^- [ a_m^2 a_n \sin(2\alpha_m - \alpha_n) \\
& + b_m^2 b_n \sin(2\beta_m - \beta_n) ] \cosh k_{2mn}^- Z \\
& + F_{n2m}^+ [ a_n b_m^2 \sin(\alpha_n + 2\beta_m) \\
& + b_n a_m^2 \sin(\beta_n + 2\alpha_m) ] \cosh k_{n2m}^+ Z \\
& + F_{n2m}^- [ a_n b_m^2 \sin(\alpha_n - 2\beta_m) \\
& + b_n a_m^2 \sin(\beta_n - 2\alpha_m) ] \cosh k_{n2m}^- Z \\
& + F_{2nm}^+ [ a_n^2 b_m \sin(2\alpha_n + \beta_m) \\
& + b_n^2 a_m \sin(2\beta_n + \alpha_m) ] \cosh k_{2nm}^+ Z \\
& + F_{2nm}^- [ a_n^2 b_m \sin(2\alpha_n - \beta_m) \\
& + b_n^2 a_m \sin(2\beta_n - \alpha_m) ] \cosh k_{2nm}^- Z \\
& + F_{mnn}^+ [ a_m a_n b_m \sin(\alpha_m + \alpha_n - \beta_m) \\
& + b_m b_n a_m \sin(\beta_m + \beta_n - \alpha_m) ] \cosh k_{mnn}^+ Z \\
& + F_{mnn}^- [ a_m a_n b_m \sin(\alpha_m - \alpha_n - \beta_m) \\
& + b_m b_n a_m \sin(\beta_m - \beta_n - \alpha_m) ] \cosh k_{mnn}^- Z \\
& + F_{mnn}^+ [ a_m a_n b_n \sin(\alpha_m + \alpha_n - \beta_n) \\
& + b_m b_n a_n \sin(\beta_m + \beta_n - \alpha_n) ] \cosh k_{mnn}^+ Z \\
& + F_{mnn}^- [ a_m a_n b_n \sin(\alpha_m - \alpha_n + \beta_n) \\
& + b_m b_n a_n \sin(\beta_m - \beta_n + \alpha_n) ] \cosh k_{mnn}^- Z \\
& + F_{mnn}^+ [ a_m b_m b_n \sin(\alpha_m + \beta_m + \beta_n) \\
& + a_m b_m a_n \sin(\alpha_m + \beta_m + \alpha_n) ] \cosh k_{mnn}^+ Z \\
& + F_{mnn}^- [ a_m b_m b_n \sin(\alpha_m + \beta_m - \beta_n) \\
& + a_m b_m a_n \sin(\alpha_m + \beta_m - \alpha_n) ] \cosh k_{mnn}^- Z \\
& + F_{nmm}^+ [ a_n b_m b_n \sin(\alpha_n + \beta_m + \beta_n) \\
& + a_n a_m b_n \sin(\alpha_n + \alpha_m + \beta_n) ] \cosh k_{nmm}^+ Z \\
& + F_{nmm}^- [ a_n b_m b_n \sin(\alpha_n - \beta_m + \beta_n) \\
& + a_n a_m b_n \sin(\alpha_n - \alpha_m + \beta_n) ] \cosh k_{nmm}^- Z, (53)
\end{aligned}$$

$$\begin{aligned}
\omega_{21} = & - \frac{\omega_0 a_m^2}{2\omega_m} \Gamma_{36}(F_{2m}, G_{2m}, k_{2m}, k_m, \omega_m) \\
& - \frac{\omega_0 b_m^2}{2\omega_m} \Gamma_{37}(F_{mm}^+, G_{mm}^+, G_{mm}^-, k_{mm}^+, \\
& k_m, \omega_m, \gamma_m^-, m_1) \\
& - \frac{\omega_0 a_n^2}{2\omega_m} \Gamma_{38}(F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, \\
& k_{mn}^+, k_{mn}^-, k_n, k_m, \omega_n, \omega_m, \gamma_{mn}^+, \gamma_n^+) - \\
& \frac{\omega_0 b_n^2}{2\omega_m} \Gamma_{38}(F_{nm}^+, -F_{nm}^-, G_{nm}^+, G_{nm}^-, \\
& k_{nm}^+, k_{nm}^-, k_n, k_m, \omega_n, \omega_m, \gamma_{nm}^-, \gamma_n^+), (54)
\end{aligned}$$

$$\begin{aligned}
\omega_{22} = & - \frac{\omega_0 a_m^2}{2\omega_m} \Gamma_{37}(F_{mm}^+, G_{mm}^+, G_{mm}^-, \\
& k_{mm}^+, k_m, \omega_m, \gamma_m^-, m_1) \\
& - \frac{\omega_0 b_m^2}{2\omega_m} \Gamma_{36}(F_{2m}, G_{2m}, k_{2m}, k_m, \omega_m) \\
& - \frac{\omega_0 a_n^2}{2\omega_m} \Gamma_{38}(F_{nm}^+, -F_{nm}^-, G_{nm}^+, G_{nm}^-, \\
& k_{nm}^+, k_{nm}^-, k_n, k_m, \omega_n, \omega_m, \gamma_{nm}^-, \gamma_n^+) \\
& - \frac{\omega_0 b_n^2}{2\omega_m} \Gamma_{38}(F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, \\
& k_{mn}^+, k_{mn}^-, k_n, k_m, \omega_n, \omega_m, \gamma_{mn}^+, \gamma_n^+), (55)
\end{aligned}$$

$$\begin{aligned}
\omega_{23} = & - \frac{\omega_0 a_m^2}{2\omega_n} \Gamma_{38}(F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, \\
& k_{mn}^+, k_{mn}^-, k_m, k_n, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_m^+) \\
& - \frac{\omega_0 b_m^2}{2\omega_n} \Gamma_{38}(F_{nm}^+, -F_{nm}^-, G_{nm}^+, G_{nm}^-, \\
& k_{nm}^+, k_{nm}^-, k_m, k_n, \omega_m, \omega_n, \gamma_{nm}^-, \gamma_m^+) \\
& - \frac{\omega_0 a_n^2}{2\omega_n} \Gamma_{36}(F_{2n}, G_{2n}, k_{2n}, k_n, \omega_n) \\
& - \frac{\omega_0 b_n^2}{2\omega_n} \Gamma_{37}(F_{nn}^+, G_{nn}^+, G_{nn}^-, \\
& k_{nn}^+, k_n, \omega_n, \gamma_n^-, n_1), (56)
\end{aligned}$$

$$\begin{aligned}
\omega_{24} = & - \frac{\omega_0 a_m^2}{2\omega_n} \Gamma_{38}(F_{nm}^+, -F_{nm}^-, G_{nm}^+, G_{nm}^-, \\
& k_{nm}^+, k_{nm}^-, k_m, k_n, \omega_m, \omega_n, \gamma_{nm}^-, \gamma_m^+) \\
& - \frac{\omega_0 b_m^2}{2\omega_n} \Gamma_{38}(F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, \\
& k_{mn}^+, k_{mn}^-, k_m, k_n, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_m^+) \\
& - \frac{\omega_0 a_n^2}{2\omega_n} \Gamma_{37}(F_{nn}^+, G_{nn}^+, G_{nn}^-, \\
& k_{nn}^+, k_n, \omega_n, \gamma_n^-, n_1) \\
& - \frac{\omega_0 b_n^2}{2\omega_n} \Gamma_{36}(F_{2n}, G_{2n}, k_{2n}, k_n, \omega_n), (57)
\end{aligned}$$

其中,各传递函数和有关的波数如下(同样,对于具有同等表达形式的两个对称项,仅列出涉及下标 m 的项):

$$\begin{aligned}
G_{m0} = & \frac{\omega_2}{\omega_0} + a_m^2 \Lambda_{31}(F_{2m}, G_{2m}, k_{2m}, k_m, \omega_m) \\
& + b_m^2 \Lambda_{32}(F_{mm}^+, G_{mm}^+, G_{mm}^-, k_{mm}^+, k_m, \omega_m, m_1) \\
& + a_n^2 \Lambda_{33}(F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, \\
& k_{mn}^+, k_{mn}^-, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_m^+, \gamma_n^+)
\end{aligned}$$

$$+ b_n^2 A_{33} (F_{nm}^+, -F_{nm}^-, G_{nm}^+, G_{nm}^-, k_{nm}^+, k_{nm}^-, \omega_m, \omega_n, \gamma_{nm}^-, \gamma_m^+, \gamma_n^+), \quad (58)$$

$$G_{0m} = \frac{\omega_2}{\omega_0} + a_m^2 A_{32} (F_{mm}^+, G_{mm}^+, G_{mm}^-, k_{mm}^+, k_m, \omega_m, m_1) + b_m^2 A_{31} (F_{2m}, G_{2m}, k_{2m}, k_m, \omega_m) + a_n^2 A_{33} (F_{nm}^+, -F_{nm}^-, G_{nm}^+, G_{nm}^-, k_{nm}^+, k_{nm}^-, \omega_m, \omega_n, \gamma_{nm}^-, \gamma_m^+, \gamma_n^+) + b_n^2 A_{33} (F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, k_{mn}^+, k_{mn}^-, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_m^+, \gamma_n^+), \quad (59)$$

$$G_{3m} = \frac{3}{2} F_{2m} k_{2m} \omega_m \sinh k_{2m} h - \frac{1}{4\omega_m} F_{2m} k_{2m}^2 \cosh k_{2m} h + 3\omega_m F_{3m} \cosh 3k_m h + \frac{1}{2} G_{2m} \omega_m^2 + \frac{1}{8} \gamma_m^+, \quad (60)$$

$$G_{m2m}^- = \frac{1}{2} F_{2m} k_{2m} \omega_m \sinh k_{2m} h - \frac{\gamma_m^-}{\omega_m} F_{2m} \cosh k_{2m} h - \omega_m F_{m2m}^- \cosh k_{m2m}^- h + \frac{1}{2} (G_{2m} + G_{mm}^-) \omega_m^2 + \frac{1}{8} k_m^2 (3 - 8m_1^2), \quad (61)$$

$$G_{2mm}^+ = \frac{3}{2} F_{2m} k_{2m} \omega_m \sinh k_{2m} h - \frac{\gamma_m^-}{\omega_m} F_{2m} \cosh k_{2m} h + \frac{3}{2} F_{mm}^+ k_{mm}^+ \omega_m \sinh k_{mm}^+ h - \frac{(k_{mm}^+)^2}{4\omega_m} F_{mm}^+ \cosh k_{mm}^+ h + 3\omega_m F_{2mm}^+ \cosh k_{2mm}^+ h + \frac{1}{2} (G_{2m} + G_{mm}^+) \omega_m^2 + \frac{1}{8} k_m^2 (11 - 8m_1^2), \quad (62)$$

$$G_{m2n}^\pm = A_{34} (\pm F_{2n}, F_{mn}^\pm, G_{2n}, G_{mn}^\pm, \pm k_{2n}, k_{mn}^\pm, k_m, \pm k_n, \omega_m, \pm \omega_n, \pm \gamma_{mn}^+, \gamma_n^+) + (\omega_m \pm 2\omega_n) F_{m2n}^\pm \cosh k_{m2n}^\pm h, \quad (63)$$

$$G_{2mn}^\pm = A_{34} (F_{2m}, F_{mn}^\pm, G_{2m}, G_{mn}^\pm, k_{2m}, k_{mn}^\pm, \pm k_n, k_m, \pm \omega_n, \omega_m, \pm \gamma_{mn}^+, \gamma_m^+) + (2\omega_m \pm \omega_n) F_{2mn}^\pm \cosh k_{2mn}^\pm h, \quad (64)$$

$$G_{n2m}^\pm = A_{34} (F_{2m}, \pm F_{nm}^\pm, G_{2m}, G_{nm}^\pm, k_{2m}, k_{nm}^\pm,$$

$$\pm k_n, k_m, \pm \omega_n, \omega_m, \pm \gamma_{nm}^-, \gamma_m^+) + (\omega_n \pm 2\omega_m) F_{n2m}^\pm \cosh k_{n2m}^\pm h, \quad (65)$$

$$G_{2nm}^\pm = A_{34} (\pm F_{2n}, \pm F_{nm}^\pm, G_{2n}, G_{nm}^\pm, \pm k_{2n}, k_{nm}^\pm, k_m, \pm k_n, \omega_m, \pm \omega_n, \pm \gamma_{nm}^-, \gamma_n^+) + (2\omega_n \pm \omega_m) F_{2nm}^\pm \cosh k_{2nm}^\pm h, \quad (66)$$

$$G_{mnn}^\pm = A_{35} (F_{mn}^\pm, \mp F_{nm}^\mp, G_{mn}^\pm, G_{nn}^\pm, G_{nm}^\mp, k_{mn}^\pm, k_{nm}^\mp, \pm k_n, \omega_m, \pm \omega_n, \pm \gamma_{mn}^+, \pm \gamma_{nn}^-, \pm \gamma_{11}, \gamma_m^-) \pm \omega_n F_{mnn}^\pm \cosh k_{mnn}^\pm h, \quad (67)$$

$$G_{mnn}^\pm = A_{35} (F_{mn}^\pm, \mp F_{nm}^\mp, G_{nn}^-, G_{mn}^\pm, G_{nn}^\mp, k_{mn}^\pm, k_{nm}^\mp, k_m, \pm \omega_n, \omega_m, \pm \gamma_{nm}^+, \pm \gamma_{nn}^-, \pm \gamma_{11}, \gamma_n^-) + \omega_m F_{mnn}^\pm \cosh k_{mnn}^\pm h, \quad (68)$$

$$G_{mnn}^\pm = A_{36} (F_{mm}^+, F_{mn}^\pm, \pm F_{nm}^\pm, G_{mm}^+, G_{mn}^\pm, G_{nm}^\pm, k_{mm}^+, k_{mn}^\pm, k_{nm}^\pm, k_m, \pm k_n, \omega_m, \pm \omega_n, \pm \gamma_{mn}^+, \pm \gamma_{nn}^-, \pm \gamma_{11}, \gamma_m^-) + (2\omega_m \pm \omega_n) F_{mnn}^\pm \cosh k_{mnn}^\pm h, \quad (69)$$

$$G_{mnn}^\pm = A_{36} (\pm F_{nn}^+, F_{mn}^\pm, \pm F_{nm}^\pm, G_{nn}^+, G_{mn}^\pm, G_{nm}^\pm, \pm k_{nn}^+, k_{mn}^\pm, k_{nm}^\pm, \pm k_n, k_m, \pm \omega_n, \omega_m, \pm \gamma_{nm}^+, \pm \gamma_{nn}^-, \pm \gamma_{11}, \gamma_n^-) + (2\omega_n \pm \omega_m) F_{mnn}^\pm \cosh k_{mnn}^\pm h, \quad (70)$$

$$F_{3m} = \frac{A_{3m}}{3k_m \sinh 3k_m h - 9\omega_m^2 \cosh 3k_m h}, \quad (71)$$

$$F_{m2m}^- = \frac{A_{m2m}^-}{k_{m2m}^- \sinh k_{m2m}^- h - \omega_m^2 \cosh k_{m2m}^- h}, \quad (72)$$

$$F_{2mm}^+ = \frac{A_{2mm}^+}{k_{2mm}^+ \sinh k_{2mm}^+ h - 9\omega_m^2 \cosh k_{2mm}^+ h}, \quad (73)$$

$$F_{m2n}^\pm = \frac{A_{m2n}^\pm}{k_{m2n}^\pm \sinh k_{m2n}^\pm h - (\omega_m \pm 2\omega_n)^2 \cosh k_{m2n}^\pm h}, \quad (74)$$

$$F_{2mn}^\pm = \frac{A_{2mn}^\pm}{k_{2mn}^\pm \sinh k_{2mn}^\pm h - (2\omega_m \pm \omega_n)^2 \cosh k_{2mn}^\pm h}, \quad (75)$$

$$F_{n2m}^\pm = \frac{A_{n2m}^\pm}{k_{n2m}^\pm \sinh k_{n2m}^\pm h - (\omega_n \pm 2\omega_m)^2 \cosh k_{n2m}^\pm h}, \quad (76)$$

$$F_{2nm}^\pm = \frac{A_{2nm}^\pm}{k_{2nm}^\pm \sinh k_{2nm}^\pm h - (2\omega_n \pm \omega_m)^2 \cosh k_{2nm}^\pm h}, \quad (77)$$

$$F_{mnn}^\pm = \frac{A_{mnn}^\pm}{k_{mnn}^\pm \sinh k_{mnn}^\pm h - \omega_n^2 \cosh k_{mnn}^\pm h}, \quad (78)$$

$$F_{mnn}^\pm = \frac{A_{mnn}^\pm}{k_{mnn}^\pm \sinh k_{mnn}^\pm h - \omega_m^2 \cosh k_{mnn}^\pm h}, \quad (79)$$

$$F_{mnn}^{\pm} = \frac{A_{mnn}^{\pm}}{k_{mnn}^{\pm} \sinh k_{mnn}^{\pm} h - (2\omega_m \pm \omega_n)^2 \cosh k_{mnn}^{\pm} h}, \quad (80)$$

$$F_{nmn}^{\pm} = \frac{A_{nmn}^{\pm}}{k_{nmn}^{\pm} \sinh k_{nmn}^{\pm} h - (\omega_m \pm 2\omega_n)^2 \cosh k_{nmn}^{\pm} h}, \quad (81)$$

$$k_{m2m}^{-} = \sqrt{(k_m m_1)^2 + (3k_m m_2)^2}, \quad (82)$$

$$k_{2mm}^{+} = \sqrt{(3k_m m_1)^2 + (k_m m_2)^2}, \quad (83)$$

$$k_{m2n}^{\pm} = \sqrt{(k_m m_1 \pm 2k_n n_1)^2 + (k_m m_2 \pm 2k_n n_2)^2}, \quad (84)$$

$$k_{2mn}^{\pm} = \sqrt{(2k_m m_1 \pm k_n n_1)^2 + (2k_m m_2 \pm k_n n_2)^2}, \quad (85)$$

$$k_{n2m}^{\pm} = \sqrt{(\pm k_n n_1 + 2k_m m_1)^2 + (\pm k_n n_2 - 2k_m m_2)^2}, \quad (86)$$

$$k_{2nm}^{\pm} = \sqrt{(\pm 2k_n n_1 + k_m m_1)^2 + (\pm 2k_n n_2 - k_m m_2)^2}, \quad (87)$$

$$k_{mnn}^{\pm} = \sqrt{(k_n n_1)^2 + (2k_m m_2 \pm k_n n_2)^2}, \quad (88)$$

$$k_{mnn}^{\pm} = \sqrt{(k_m m_1)^2 + (k_m m_2 \pm 2k_n n_2)^2}, \quad (89)$$

$$k_{mnn}^{\pm} = \sqrt{(2k_m m_1 \pm k_n n_1)^2 + (k_n n_2)^2}, \quad (90)$$

$$k_{nmn}^{\pm} = \sqrt{(k_m m_1 \pm 2k_n n_1)^2 + (k_m m_2)^2}. \quad (91)$$

$$\begin{aligned} & \Gamma_{36}(F_{2m}, G_{2m}, k_{2m}, k_m, \omega_m) \\ &= \omega_m^2 F_{2m} k_{2m} \sinh k_{2m} h - k_{2m}^2 F_{2m} \cosh k_{2m} h \\ &+ \frac{k_m^4 - 5k_m^2 \omega_m^4}{4\omega_m^3} + \frac{G_{2m}(k_m^2 - \omega_m^4)}{2\omega_m}, \quad (92) \end{aligned}$$

$$\begin{aligned} & \Gamma_{37}(F_{mn}^+, G_{mn}^+, G_{mn}^-, k_{mn}^+, k_m, \omega_m, \gamma_m^-, m_1) \\ &= \omega_m^2 F_{mn}^+ k_{mn}^+ \sinh k_{mn}^+ h \\ &- (k_{mn}^+)^2 F_{mn}^+ \cosh k_{mn}^+ h + \frac{(\gamma_m^-)^2}{2\omega_m^3} \\ &+ \frac{1}{2} k_m^2 \omega_m (3 - 8m_1^2) \\ &+ \frac{(G_{mn}^+ - G_{mn}^-)(k_m^2 - \omega_m^4)}{2\omega_m}, \quad (93) \end{aligned}$$

$$\begin{aligned} & \Gamma_{38}(F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, k_{mn}^+, \\ & k_{mn}^-, k_m, k_n, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_n^+) \\ &= \Gamma_{32}^-(F_{mn}^+, k_{mn}^+, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_n^+) \\ &+ \Gamma_{32}^-(F_{mn}^-, k_{mn}^-, \omega_m, -\omega_n, -\gamma_{mn}^+, \gamma_n^+) \\ &+ \frac{(\gamma_{mn}^+)^2}{2\omega_m \omega_n^2} - \frac{\gamma_{mn}^+ (\omega_m^2 + \omega_n^2)}{\omega_n} \\ &+ \frac{(G_{mn}^+ - G_{mn}^-)(k_n^2 - \omega_n^4)}{2\omega_n} \\ &+ \frac{k_m^2 \omega_n^2 - 2k_n^2 \omega_m^2}{2\omega_m}, \quad (94) \end{aligned}$$

其中,所引入函数的具体表达式或新引入的函数为

$$\begin{aligned} & \Lambda_{31}(F_{2m}, G_{2m}, k_{2m}, k_m, \omega_m) \\ &= \frac{1}{2} F_{2m} k_{2m} \omega_m \sinh k_{2m} h \\ &- \frac{1}{4\omega_m} F_{2m} k_{2m}^2 \cosh k_{2m} h \\ &+ \frac{1}{2} G_{2m} \omega_m^2 - \frac{5}{8} k_m^2, \quad (95) \end{aligned}$$

$$\begin{aligned} & \Lambda_{32}(F_{mn}^+, G_{mn}^+, G_{mn}^-, k_{mn}^+, k_m, \omega_m, m_1) \\ &= \frac{1}{2} F_{mn}^+ k_{mn}^+ \omega_m \sinh k_{mn}^+ h \\ &- \frac{1}{4\omega_m} F_{mn}^+ (k_{mn}^+)^2 \cosh k_{mn}^+ h \\ &+ \frac{1}{2} (G_{mn}^+ + G_{mn}^-) \omega_m^2 + \frac{1}{4} k_m^2 (3 - 8m_1^2), \quad (96) \end{aligned}$$

$$\begin{aligned} & \Lambda_{33}(F_{mn}^+, F_{mn}^-, G_{mn}^+, G_{mn}^-, k_{mn}^+, \\ & k_{mn}^-, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_m^+, \gamma_n^+) \\ &= \frac{1}{2} F_{mn}^+ k_{mn}^+ \omega_m \sinh k_{mn}^+ h - \frac{\gamma_{mn}^+ + \gamma_n^+}{2\omega_n} F_{mn}^+ \cosh k_{mn}^+ h \\ &+ \frac{1}{2} F_{mn}^- k_{mn}^- \omega_m \sinh k_{mn}^- h - \frac{\gamma_{mn}^+ - \gamma_n^+}{2\omega_n} F_{mn}^- \cosh k_{mn}^- h \\ &+ \frac{1}{2} (G_{mn}^+ + G_{mn}^-) \omega_n^2 - \frac{\gamma_m^+}{2} \left( \frac{\omega_m}{\omega_n} + \frac{\omega_n}{\omega_m} \right) \\ &+ \frac{1}{4} (\gamma_m^+ - 2\gamma_n^+), \quad (97) \end{aligned}$$

$$\begin{aligned} & \Lambda_{34}(F_{2n}, F_{mn}^{\pm}, G_{2n}, G_{mn}^{\pm}, k_{2n}, \\ & k_{mn}^{\pm}, k_m, k_n, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_n^+) \\ &= \frac{(2\omega_n + \omega_m)}{2} F_{2n} k_{2n} \sinh k_{2n} h - \frac{\gamma_{mn}^+}{\omega_m} F_{2n} \cosh k_{2n} h \\ &+ \frac{(\omega_m + 2\omega_n)}{2} F_{mn}^{\pm} k_{mn}^{\pm} \sinh k_{mn}^{\pm} h \\ &- \frac{\gamma_{mn}^+ + \gamma_n^+}{2\omega_n} F_{mn}^{\pm} \cosh k_{mn}^{\pm} h \\ &+ \frac{1}{2} (G_{2n} \omega_m^2 + G_{mn}^{\pm} \omega_n^2) + \frac{1}{8} (k_m^2 + 2k_n^2) \\ &- \frac{\gamma_{mn}^+}{4} \left( \frac{\omega_m}{\omega_n} + \frac{\omega_n}{\omega_m} \right) + \frac{1}{4} \left( \frac{k_n^2 \omega_m}{\omega_n} + \frac{k_m^2 \omega_n}{\omega_m} \right), \quad (98) \end{aligned}$$

$$\begin{aligned} & \Lambda_{35}(F_{mn}^{\pm}, \mp F_{nm}^{\mp}, G_{nm}^-, G_{mn}^{\pm}, G_{nm}^{\mp}, k_{mn}^{\pm}, \\ & k_{nm}^{\mp}, k_n, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_{nm}^-, \gamma_{11}, \gamma_m^-) \\ &= \frac{1}{2} F_{mn}^{\pm} k_{mn}^{\pm} \omega_n \sinh k_{mn}^{\pm} h - \frac{\gamma_m^- + \gamma_{nm}^-}{2\omega_n} F_{mn}^{\pm} \cosh k_{mn}^{\pm} h \end{aligned}$$

$$\begin{aligned}
& -\frac{1}{2}F_{nm}^{\mp}k_{nm}^{\mp}\omega_n\sinh k_{nm}^{\mp}h - \frac{\gamma_m^- - \gamma_{mn}^+}{2\omega_m}F_{nm}^{\mp}\cosh k_{nm}^{\mp}h \\
& + \frac{1}{2}G_{mm}^-\omega_n^2 + \frac{1}{2}(G_{mn}^+ + G_{nm}^{\mp})\omega_m^2 \\
& + \frac{1}{4}(k_n^2 - 2\gamma_m^-) - \frac{\gamma_{11}}{2}\left(\frac{\omega_m}{\omega_n} + \frac{\omega_n}{\omega_m}\right), \quad (99)
\end{aligned}$$

$$\begin{aligned}
& A_{36}(F_{mm}^+, F_{mn}^{\pm}, \pm F_{nm}^{\pm}, G_{mm}^+, G_{mn}^{\pm}, G_{nm}^{\pm}, k_{mm}^+, \\
& k_{mn}^{\pm}, k_{nm}^{\pm}, k_m, \pm k_n, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_{mn}^-, \gamma_{11}, \gamma_m^-) \\
& = \frac{2\omega_m + \omega_n}{2}F_{mm}^+k_{mm}^+\sinh k_{mm}^+h - \frac{\gamma_{11}}{\omega_n}F_{mm}^+\cosh k_{mm}^+h \\
& + \frac{2\omega_m + \omega_n}{2}F_{mn}^{\pm}k_{mn}^{\pm}\sinh k_{mn}^{\pm}h \\
& - \frac{\gamma_m^- + \gamma_{mn}^-}{2\omega_m}F_{mn}^{\pm}\cosh k_{mn}^{\pm}h \\
& + \frac{2\omega_m + \omega_n}{2}F_{nm}^{\pm}k_{nm}^{\pm}\sinh k_{nm}^{\pm}h \\
& - \frac{\gamma_m^- + \gamma_{mn}^+}{2\omega_m}F_{nm}^{\pm}\cosh k_{nm}^{\pm}h + \frac{1}{2}G_{mm}^+\omega_n^2 \\
& + \frac{1}{2}(G_{mn}^{\pm} + G_{nm}^{\pm})\omega_m^2 + \frac{1}{4}(k_n^2 + 2k_m^2) \\
& - \frac{1}{2}(\gamma_m^- - k_m^2) - \frac{\gamma_{11}}{2}\left(\frac{\omega_m}{\omega_n} + \frac{\omega_n}{\omega_m}\right) \\
& + \frac{\omega_n}{\omega_m} + \frac{1}{2}\left(\frac{k_n^2\omega_m}{\omega_n} + \frac{k_m^2\omega_n}{\omega_m}\right), \quad (100)
\end{aligned}$$

$$\begin{aligned}
A_{3m} & = 5\omega_m^2 F_{2m} k_{2m} \sinh k_{2m} h \\
& - 2k_{2m}^2 F_{2m} \cosh k_{2m} h \\
& + \frac{k_m^4 - k_m^2 \omega_m^4}{4\omega_m^3} - \frac{G_{2m}(k_m^2 - \omega_m^4)}{2\omega_m}, \quad (101)
\end{aligned}$$

$$\begin{aligned}
A_{m2m}^- & = -\omega_m^2 F_{2m} k_{2m} \sinh k_{2m} h \\
& + k_{2m}^2 m_2^2 F_{2m} \cosh k_{2m} h \\
& + \frac{(4m_2^4 - 1)k_m^4 + (5 - 8m_2^2)k_m^2 \omega_m^4}{4\omega_m^3} \\
& - \frac{(G_{2m} - G_{mm}^-)(k_m^2 - \omega_m^4)}{2\omega_m}, \quad (102)
\end{aligned}$$

$$\begin{aligned}
A_{2mm}^+ & = 5\omega_m^2 F_{2m} k_{2m} \sinh k_{2m} h \\
& + (3m_2^2 - 2)k_{2m}^2 F_{2m} \cosh k_{2m} h \\
& - \frac{(G_{2m} + G_{mm}^+)(k_m^2 - \omega_m^4)}{2\omega_m} \\
& + 5\omega_m^2 F_{mm}^+ k_{mm}^+ \sinh k_{mm}^+ h - 2(k_{mm}^+)^2 F_{mm}^+ \cosh k_{mm}^+ h \\
& + \frac{(4m_1^4 - 1)k_m^4 + 3(8m_2^2 - 1)k_m^2 \omega_m^4}{4\omega_m^3}, \quad (103)
\end{aligned}$$

$$\begin{aligned}
A_{m2n}^{\pm} & = \Gamma_{31}(F_{2n}, k_{2n}, \pm \omega_n, \omega_m, \pm \gamma_{nm}^+) \\
& + \Gamma_{32}^+(F_{mn}^{\pm}, k_{mn}^{\pm}, \pm \omega_n, \omega_m, \pm \gamma_{mn}^+, \gamma_n^+) \\
& + \Gamma_{33}^+(G_{2n}, G_{mn}^{\pm}, \pm k_n, k_m, \pm \omega_n, \omega_m, \pm \gamma_{mn}^+), \quad (104)
\end{aligned}$$

$$\begin{aligned}
A_{2mn}^{\pm} & = \Gamma_{31}(F_{2m}, k_{2m}, \omega_m, \pm \omega_n, \pm \gamma_{mn}^+) \\
& + \Gamma_{32}^+(F_{mn}^{\pm}, k_{mn}^{\pm}, \omega_m, \pm \omega_n, \pm \gamma_{mn}^+, \gamma_n^+) \\
& + \Gamma_{33}^-(G_{2m}, G_{mn}^{\pm}, k_m, \pm k_n, \omega_m, \pm \omega_n, \pm \gamma_{mn}^+), \quad (105)
\end{aligned}$$

$$\begin{aligned}
A_{n2m}^{\pm} & = \pm \Gamma_{31}(F_{2m}, k_{2m}, \omega_m, \pm \omega_n, \pm \gamma_{nm}^-) \\
& \pm \Gamma_{32}^+(\pm F_{nm}^{\pm}, k_{nm}^{\pm}, \omega_m, \pm \omega_n, \pm \gamma_{nm}^-, \gamma_m^+) \\
& \pm \Gamma_{33}^-(G_{2m}, G_{nm}^{\pm}, k_m, \pm k_n, \omega_m, \pm \omega_n, \pm \gamma_{nm}^-), \quad (106)
\end{aligned}$$

$$\begin{aligned}
A_{2nm}^{\pm} & = \pm \Gamma_{31}(F_{2n}, k_{2n}, \pm \omega_n, \omega_m, \pm \gamma_{nm}^-) \\
& \pm \Gamma_{32}^+(F_{mn}^{\pm}, \pm k_{mn}^{\pm}, \pm \omega_n, \omega_m, \pm \gamma_{nm}^-, \gamma_n^+) \pm \\
& \Gamma_{33}^-(G_{2n}, G_{nm}^{\pm}, \pm k_n, k_m, \pm \omega_n, \omega_m, \pm \gamma_{nm}^-), \quad (107)
\end{aligned}$$

$$\begin{aligned}
A_{mnm}^{\pm} & = \Gamma_{32}^-(F_{mn}^{\pm}, k_{mn}^{\pm}, \omega_m, \pm \omega_n, \pm \gamma_{mn}^+, \gamma_m^-) \\
& - \Gamma_{32}^-(\mp F_{nm}^{\mp}, k_{nm}^{\mp}, \omega_m, \mp \omega_n, \mp \gamma_{nm}^+, \gamma_m^-) \\
& + \Gamma_{34}(G_{mn}^-, G_{mn}^{\pm}, G_{nm}^{\mp}, k_m, \pm k_n, \omega_m, \pm \omega_n, \\
& \pm \gamma_{11}, \pm \gamma_{22}, \pm \gamma_{mn}^+, \pm \gamma_{nm}^-, \gamma_m^-), \quad (108)
\end{aligned}$$

$$\begin{aligned}
A_{mnn}^{\pm} & = \Gamma_{32}^-(F_{mn}^{\pm}, k_{mn}^{\pm}, \pm \omega_n, \omega_m, \pm \gamma_{mn}^+, \gamma_n^-) \\
& + \Gamma_{32}^-(\mp F_{nm}^{\mp}, k_{nm}^{\mp}, \mp \omega_n, \omega_m, \mp \gamma_{nm}^+, \gamma_n^-) \\
& + \Gamma_{34}(G_{nn}^-, G_{nn}^{\pm}, G_{nm}^{\mp}, \pm k_n, k_m, \pm \omega_n, \omega_m, \\
& \pm \gamma_{11}, \pm \gamma_{22}, \pm \gamma_{mn}^+, \pm \gamma_{nm}^-, \gamma_n^-), \quad (109)
\end{aligned}$$

$$\begin{aligned}
A_{mnn}^{\pm} & = \Gamma_{31}(F_{mm}^+, k_{mm}^+, \omega_m, \pm \omega_n, \pm \gamma_{11}) \\
& + \Gamma_{32}^+(F_{mn}^{\pm}, k_{mn}^{\pm}, \omega_m, \pm \omega_n, \pm \gamma_{nm}^-, \gamma_m^-) \\
& + \Gamma_{32}^+(\pm F_{nm}^{\pm}, k_{nm}^{\pm}, \omega_m, \pm \omega_n, \pm \gamma_{nm}^+, \gamma_m^-) \\
& + \Gamma_{35}(G_{mm}^+, G_{mn}^{\pm}, G_{nm}^{\pm}, k_m, \pm k_n, \omega_m, \pm \omega_n, \\
& \pm \gamma_{11}, \pm \gamma_{mn}^+, \pm \gamma_{nm}^-, \gamma_m^-), \quad (110)
\end{aligned}$$

$$\begin{aligned}
A_{nmm}^{\pm} & = \pm \Gamma_{31}(F_{nn}^+, k_{nn}^+, \pm \omega_n, \omega_m, \pm \gamma_{11}) \\
& \pm \Gamma_{32}^+(F_{mn}^{\pm}, k_{mn}^{\pm}, \pm \omega_n, \omega_m, \pm \gamma_{mn}^-, \gamma_n^-) \\
& \pm \Gamma_{32}^+(\pm F_{nm}^{\pm}, k_{nm}^{\pm}, \pm \omega_n, \omega_m, \pm \gamma_{nm}^+, \gamma_n^-) \\
& \pm \Gamma_{35}(G_{nn}^+, G_{mn}^{\pm}, G_{nm}^{\pm}, \pm k_n, k_m, \pm \omega_n, \omega_m, \\
& \pm \gamma_{11}, \pm \gamma_{nm}^+, \pm \gamma_{nn}^-, \gamma_n^-), \quad (111)
\end{aligned}$$

其中,新引入函数的具体表达形式为

$$\begin{aligned}
& \Gamma_{31}(F_{2m}, k_{2m}, \omega_m, \omega_n, \gamma_{mn}^+) \\
& = (\omega_n^2 + 2\omega_m \omega_n + 2\omega_m^2) F_{2m} k_{2m} \sinh k_{2m} h
\end{aligned}$$



$$\begin{aligned}
& - \left( 2\gamma_{mn}^+ \frac{\omega_n + 2\omega_m}{\omega_n} + \frac{1}{2}k_{2m}^2 \right) F_{2m} \cosh k_{2m} h, \quad (112) \\
& \Gamma_{32}^\pm (F_{mn}^+, k_{mn}^+, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_m^+) \\
= & \left[ \omega_m^2 + \frac{1}{2}(\omega_m + \omega_n)^2 \pm \omega_m(\omega_m + \omega_n) \right] \\
& \times F_{mn}^+ k_{mn}^+ \sinh k_{mn}^+ h \\
& - \left[ (\gamma_{mn}^+ + \gamma_m^+) \left( \frac{\omega_n + \omega_m}{\omega_m} \pm 1 \right) \right. \\
& \left. + \frac{1}{2}(k_{mn}^+)^2 \right] F_{mn}^+ \cosh k_{mn}^+ h, \quad (113)
\end{aligned}$$

$$\begin{aligned}
& \Gamma_{33}^\pm (G_{2m}^+, G_{mn}^+, k_m, k_n, \omega_m, \omega_n, \gamma_{mn}^+, \gamma_m^+) \\
= & - \frac{1}{2} \left( G_{2m}^+ \frac{k_n^2 - \omega_n^4}{\omega_n} + G_{mn}^+ \frac{k_m^2 - \omega_m^4}{\omega_m} \right) \\
& + \frac{\gamma_{mn}^+ (\gamma_{mn}^+ + 2k_m^2)}{4\omega_m^2 \omega_n} + \frac{k_n^2 \omega_m^2 - 2\gamma_{mn}^+ (\omega_m^2 + \omega_n^2)}{4\omega_n} \\
& \pm \frac{1}{2} [\gamma_{mn}^+ (\omega_m^2 + \omega_n^2) - (k_n^2 \omega_m^2 + k_m^2 \omega_n^2)] \\
& \times \left( \frac{1}{\omega_m} + \frac{1}{\omega_n} \right), \quad (114)
\end{aligned}$$

$$\begin{aligned}
& \Gamma_{34} (G_{mn}^-, G_{mn}^+, G_{nm}^-, k_m, k_n, \omega_m, \omega_n, \\
& \gamma_{11}, \gamma_{22}, \gamma_{mn}^+, \gamma_{mn}^-, \gamma_m^-) \\
= & (G_{mn}^+ - G_{nm}^-) \frac{k_m^2 - \omega_m^4}{2\omega_m} \\
& - G_{mn}^- \frac{k_n^2 - \omega_n^4}{2\omega_n} + \frac{\gamma_{mn}^+ \gamma_{mn}^- + 2\gamma_m^- \gamma_{22}}{2\omega_m^2 \omega_n} \\
& + \frac{k_n^2 \omega_m^2 - 2\gamma_m^- \omega_n^2}{2\omega_n} - \frac{\gamma_{11} (\omega_m^2 + \omega_n^2)}{\omega_m}, \quad (115)
\end{aligned}$$

$$\begin{aligned}
& \Gamma_{35} (G_{mn}^+, G_{mn}^+, G_{nm}^+, k_m, k_n, \omega_m, \omega_n, \gamma_{11}, \\
& \gamma_{mn}^+, \gamma_{mn}^-, \gamma_m^-) \\
= & - \frac{1}{2} \left[ G_{mn}^+ \frac{k_n^2 - \omega_n^4}{\omega_n} + (G_{mn}^+ + G_{nm}^+) \frac{k_m^2 - \omega_m^4}{\omega_m} \right] \\
& + \frac{\gamma_{mn}^+ \gamma_{mn}^- + 2\gamma_m^- \gamma_{11}}{2\omega_m^2 \omega_n} \\
& + \frac{3k_n^2 \omega_m^2 - 2(\gamma_m^- - 2k_m^2) \omega_n^2 - 4\gamma_{11} (\omega_m^2 + \omega_n^2)}{2\omega_n} \\
& + \frac{(k_n^2 \omega_m^2 + k_m^2 \omega_n^2) - \gamma_{11} (\omega_m^2 + \omega_n^2)}{\omega_m} \\
& + 4k_m^2 \omega_m m_2^2. \quad (116)
\end{aligned}$$

其中,  $\gamma_{11} = k_m k_n m_1 n_1$ ,  $\gamma_{22} = k_m k_n m_2 n_2$ .

显见,  $\zeta_3$  的前四项呈与  $\zeta_1$  类似的表达形式, 而  $\zeta_3$  的其余众多项与  $\phi_3$  的项一一对称对应, 从中可抽取出下列丰富、完整的第三阶位相组合, 表明对称

性、位相决定第三阶波: 同波同相和,  $3\alpha_m, 3\beta_m, 3\alpha_n, 3\beta_n$ ; 同波异相倍和差,  $\alpha_m - 2\beta_m, \beta_m - 2\alpha_m, \alpha_n - 2\beta_n, \beta_n - 2\alpha_n, 2\alpha_m + \beta_m, 2\beta_m + \alpha_m, 2\alpha_n + \beta_n, 2\beta_n + \alpha_n$ ; 异波同函数相倍和差,  $\alpha_m \pm 2\alpha_n, \beta_m \pm 2\beta_n, 2\alpha_m \pm \alpha_n, 2\beta_m \pm \beta_n$ ; 异波异函数相倍和差,  $\alpha_n \pm 2\beta_m, \beta_n \pm 2\alpha_m, 2\alpha_n \pm \beta_m, 2\beta_n \pm \alpha_m$ ; 同波异相差与异波相和差 ( $\alpha_m - \beta_m$ )  $\pm \alpha_n, (\beta_m - \alpha_m) \pm \beta_n, \alpha_m \pm (\alpha_n - \beta_n), \beta_m \pm (\beta_n - \alpha_n)$ ; 同波异相和与异波相和差 ( $\alpha_m + \beta_m$ )  $\pm \beta_n, (\alpha_m + \beta_m) \pm \alpha_n, (\alpha_n + \beta_n) \pm \beta_m, (\alpha_n + \beta_n) \pm \alpha_m$ .

若要论  $n$  色  $n$  向波系统, 则其某种高阶的位相组合也不难一一区分和得到. 以此刻画该波系统的基本构造.

依据上述三阶解, 即可得到其一系列运动学、动力学变量表达式<sup>[12]</sup>, 进而可最终推导出作用于直立防波堤上的三阶波浪力和波浪力矩表达式.

#### 4. 普适法则

黄虎等<sup>[13]</sup>曾从部分反射的三阶单色单向短峰波系统中推断出适用于高阶的一般法则: 倍频率通向短峰波. 现在, 若将三阶双色双向波系统发展到高阶, 其频率个数将作何变化呢? 进一步, 若将双色双向推广至  $j$  色  $j$  向, 其频率个数又是怎样的呢? 可以预先推断:  $j$  色  $j$  向波必与单色单向波一脉相承, 同样对前者可做如同后者<sup>[13]</sup>的如下频率个数的剖析.

前两阶频率为:  $\omega_0 \neq 0, \omega_1 = 0$ , 第 3 阶频率  $\omega_2$  存在 4 个值:  $\omega_{21}, \omega_{22}, \omega_{23}, \omega_{24}$ . 其原因在于:  $\omega_1, \omega_2$  是通过消除各阶方程中出现的长期项而得到的, 求解  $\omega_1$  时, 仅与 (10) 式右端项  $F_2$  中的  $-2\omega_0 \omega_1 \phi_{1u}$  项相关, 而在求解  $\omega_2$  时, 则相关于 (10) 式右端项  $F_3$  中的  $-(2\omega_0 \omega_2 + \omega_1^2) \phi_{1u}$  项和以  $\phi_1$  中的 4 项为因子的项. 如此而论, 偶数阶频率  $\omega_3, \omega_5, \dots$ , 将依次来源于  $-2\omega_0 \omega_3 \phi_{1u}, -2\omega_0 \omega_5 \phi_{1u}, \dots$ , 则  $\omega_3 = \omega_5 = 0$ ; 奇数阶频率  $\omega_4, \omega_6, \dots$ , 将依次来源于 (注意: 偶数阶频率已为零)  $-(2\omega_0 \omega_4 + \omega_2^2) \phi_{1u}, -(2\omega_0 \omega_6 + 2\omega_1 \omega_5 + 2\omega_2 \omega_4 + \omega_3^2) \phi_{1u}, \dots$ , 以及相应各阶其他以  $\phi_1$  中的 4 项为因子的项. 由于  $\omega_2, \omega_4, \omega_6, \dots$ , 各自相关于求解方程中的 4 个不同长期项, 则由消除长期项的条件, 再纳入一阶频率  $\omega_0$ , 则奇数阶频率个数构成一个等比数列:  $1, 4, 4^2, 4^3, \dots$ . 如果扩展至高阶  $j$  色  $j$  向波, 则不难推论出一项普适法则: 偶数阶频率皆为零; 奇数阶频率的个数则呈现一个等比数列:  $1,$

$2j, (2j)^2, (2j)^3, \dots$ .

## 5. 结 论

从单色单向短峰波理论发展到双色双向波, 并不仅仅是添加一色一向. 这将带来高阶推导的繁冗性、内容要旨的丰富性和形式结构的对称性. 本文

详尽刻画了双色双向波三阶解, 从中提炼出其构造要素的基本位相组合, 并意外发现了奇数阶频率个数构成一等比数列的多色多向波演变的普适法则. 这表明对称决定双色双向波相互作用恰为普适原理对称决定相互作用<sup>[14]</sup>的一个具体表现, 同时也见证、预示着真实海洋波动的异常复杂性.

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# Third-order theory for bichromatic bidirectional water waves partially reflected from vertical breakwaters<sup>\*</sup>

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## Abstract

To overcome the internal limitation of standing wave forces on vertical breakwaters and describe multichromatic multidirectional water waves in front of the breakwaters, a third-order theory, including the classical solutions of two-dimensional long-crested waves and monochromatic unidirectional short-crested waves, for bichromatic bidirectional water waves partially reflected from the breakwaters is developed. A universal law applicable to a higher-order multichromatic multidirectional wave system is proposed, that is, the even-order frequencies are zero, the numbers of the odd-order frequencies constitute a geometric progression.

**Keywords:** bichromatic bidirectional waves, a third-order theory, a universal law, partially reflected waves

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