

一类非线性演化方程新的显式行波解^{*}

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借助 Mathematica 软件, 采用三角函数法和吴文俊消元法, 获得了一类非线性演化方程 $u_{tt} + au_{xx} + bu + cu^3 = 0$ 的三组行波解, 其中包括新的行波解、扭状孤波解和钟状孤波解. 从而作为该方程的特例, 如 Duffing 方程, Klein-Gordon 方程、Landau-Ginzburg-Higgs 方程和 ϕ^4 方程等也都获得了相应的若干行波解. 这种方法也适用于其他非线性方程.

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1 引 言

随着科学技术的飞速发展, 现代科学研究的核心已逐步从线性转向非线性, 很多非线性科学问题的研究最终可用非线性方程这个数学模型来简炼而准确地描述, 而求解非线性方程长期以来一直是物理学家和数学家研究的重要课题, 虽然现在有很多方法, 如反散射法, Hopf-Cole 变换, Miura 变换, Darboux 变换和 Backlund 变换^[1-5]被有效地运用于具体的非线性方程, 但求解非线性方程仍是一个长期而艰巨的任务.

本文借助 Mathematica 软件, 采用三角函数法^[6]和吴文俊消元法, 获得了一类非线性演化方程

$$u_{tt} + au_{xx} + bu + cu^3 = 0 \quad (1)$$

(a, b, c 为常数) 的三组行波解, 从而作为该方程的特例许多著名的方程获得了相应的若干行波解.

2 非线性演化方程(1)新的行波解及孤波解

对于非线性演化方程(1), 设

$$u(x, t) = \varphi(\xi), \quad \xi = \lambda(x - kt + c_0), \quad (2)$$

其中 λ, k, c_0 为待定的常数.

将(2)式代入(1)式, 得到方程(1)对应的常微分方程

$$\lambda^2(k^2 + a) \frac{d^2\varphi}{d\xi^2} + b\varphi + c\varphi^3 = 0. \quad (3)$$

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设方程(3)有如下形式的行波解:

$$\varphi(\xi) = B_1 \sin \omega + A_1 \cos \omega + A_0, \quad \text{且 } \frac{d\omega}{d\xi} = \sin \omega, \quad (4)$$

其中 A_0, A_1, B_1 为待定常数.

将(4)式代入(3)式, 借助 Mathematica 得

$$\begin{aligned} \lambda^2(k^2 + a) \frac{d^2 \varphi}{d\xi^2} + b\varphi + c\varphi^3 = & \lambda^2(k^2 + a)(-B_1 \sin \omega + 2B_1 \sin \omega \cos^2 \omega - 2A_1 \cos \omega \\ & + 2A_1 \cos^3 \omega) + b(B_1 \sin \omega + A_1 \cos \omega + A_0) \\ & + c[(3B_1^2 A_0 + A_0^3 + 3A_1 B_1^2 \cos \omega + 3A_1 A_0^2 \cos \omega \\ & - 3A_0 B_1^2 \cos^2 \omega + 3A_1^2 A_0 \cos^2 \omega + A_1^3 \cos^3 \omega \\ & - 3A_1 B_1^2 \cos^3 \omega) + (B_1^3 + 3B_1 A_0^2 + 6A_0 A_1 B_1 \cos \omega \\ & + 3B_1 A_1^2 \cos^2 \omega - B_1^3 \cos^2 \omega) \sin \omega] = 0. \end{aligned}$$

分别令上式中的常数项 $\sin \omega, \cos \omega, \sin \omega \cos \omega, \cos^2 \omega, \sin \omega \cos^2 \omega$, 及 $\cos^3 \omega$ 的系数为零, 得

$$bA_0 + c(3B_1^2 A_0 + A_1^3) = 0, \quad (5a)$$

$$-\lambda^2(k^2 + a)B_1 + bB_1 + c(B_1^3 + 3B_1 A_0^2) = 0, \quad (5b)$$

$$-2\lambda^2(k^2 + a)A_1 + bA_1 + c(3A_1 B_1^2 + 3A_1 A_0^2) = 0, \quad (5c)$$

$$6cA_0 A_1 B_1 = 0, \quad (5d)$$

$$c(-3A_0 B_1^2 + 3A_0 A_1^2) = 0, \quad (5e)$$

$$2\lambda^2(k^2 + a)B_1 + c(-B_1^3 + 3B_1 A_1^2) = 0, \quad (5f)$$

$$2\lambda^2(k^2 + a)A_1 + c(A_1^3 - 3A_1 B_1^2) = 0. \quad (5g)$$

利用吴文俊代数消元法解上述关于 $A_1, B_1, A_0, k, \lambda$ 的超定代数方程组(5a)–(5g), 得到情况 1

$$A_0 = B_1 = 0,$$

$$A_1 = \pm \sqrt{\frac{-b}{c}}, \quad \lambda = \pm \sqrt{\frac{b}{2(k^2 + a)}}, \quad b(k^2 + a) > 0, \quad bc < 0;$$

情况 2

$$A_0 = A_1 = 0,$$

$$B_1 = \pm \sqrt{\frac{-2b}{c}}, \quad \lambda = \pm \sqrt{\frac{-b}{k^2 + a}}, \quad b(k^2 + a) < 0, \quad bc < 0;$$

情况 3

$$A_0 = 0, A_1^2 + B_1^2 = 0,$$

$$A_1 = \pm \sqrt{\frac{-b}{c}}, \quad \lambda = \pm \sqrt{\frac{2b}{k^2 + a}}, \quad b(k^2 + a) > 0, \quad bc < 0.$$

对于方程

$$\frac{d\omega}{d\xi} = \sin \omega,$$

通过分离变量并且两边积分, 得

$$\sin \omega = \frac{2 A \exp(\pm \xi)}{A^2 \exp(\pm 2 \xi) + 1} \Big|_{A=1} = \operatorname{sech} \xi, \quad (6)$$

或者

$$\cos \omega = \frac{1 - A^2 \exp(\pm 2 \xi)}{A^2 \exp(\pm 2 \xi) + 1} \Big|_{A=1} = \mp \tanh \xi, \quad (7)$$

其中 $A \neq 0$ 为积分常数.

注:下面当用到(7)式时,为方便不妨取正号.

由(4),(6),(7)及情况 1,2,3 得到

(I) 非线性演化方程(1)扭状孤波解

$$u_1(x, t) = \pm \sqrt{\frac{-b}{c}} \tanh \left[\sqrt{\frac{b}{2(k^2 + a)}} (x - kt + c_0) \right], \quad b(k^2 + a) > 0, \quad bc < 0.$$

这组解其实是文献[7]中的(13)式.

(II) 非线性演化方程(1)钟状孤波解

$$u_2(x, t) = \epsilon_1 \sqrt{\frac{-2b}{c}} \operatorname{sech} \left[\sqrt{\frac{-b}{k^2 + a}} (x - kt + c_0) \right], \quad b(k^2 + a) < 0, \quad bc < 0.$$

(III) 非线性演化方程(1)新的行波解

$$\begin{aligned} u_3(x, t) &= \pm \sqrt{\frac{-b}{c}} (\tanh \xi \pm i \operatorname{sech} \xi) = \pm \sqrt{\frac{-b}{c}} \exp(\pm i \omega) \\ &= \pm \sqrt{\frac{-b}{c}} \exp \left\{ \pm i \arcsin \left[\operatorname{sech} \left[\sqrt{\frac{2b}{k^2 + a}} (x - kt + c_0) \right] \right] \right\}, \\ &\quad b(k^2 + a) > 0, \quad bc < 0, \end{aligned}$$

或者

$$\begin{aligned} u_3(x, t) &= \pm \sqrt{\frac{-b}{c}} \exp \left\{ \pm i \arccos \epsilon_1 \tanh \left[\sqrt{\frac{2b}{k^2 + a}} (x - kt + c_0) \right] \right\}, \\ &\quad b(k^2 + a) > 0, \quad bc < 0, \end{aligned}$$

其中 $\epsilon_1 = \pm 1, i^2 = -1$.

第三组解其实是扭状孤波解和钟状孤波解的复线性组合,它是复标量场中的一种孤立波,是一种特殊的孤波解,它满足当 $|\xi| \rightarrow \infty$ 时, $u_3 \rightarrow \pm \sqrt{-b/c}$.

注:当 A_1, B_1 取实数时,第三组解不存在.

3 非线性演化方程(1)的特例及其行波解

1. 由(I),(II),(III)获得 ϕ^4 方程^[1,7,8] $u_{tt} - u_{xx} + u - u^3 = 0$ 的三组行波解

$$u_1(x, t) = \pm \tanh \left[\sqrt{\frac{1}{2(k^2 - 1)}} (x - kt + c_0) \right], \quad k^2 - 1 > 0;$$

$$u_2(x, t) = \pm \sqrt{2} \operatorname{sech} \left[\sqrt{\frac{1}{1 - k^2}} (x - kt + c_0) \right], \quad k^2 - 1 < 0;$$

$$u_3(x, t) = \pm \exp \left\{ \pm i \arcsin \operatorname{sech} \left[\sqrt{\frac{2}{k^2 - 1}} (x - kt + c_0) \right] \right\}, \quad k^2 - 1 > 0.$$

2. 由 (I), (II), (III) 获得 Klein-Gordon 方程^[4,7,9] $u_{tt} - u_{xx} + m^2 u + nu^3 = 0$ 的三组行波解

$$u_1(x, t) = \pm \frac{m}{\sqrt{-n}} \tanh \left[\sqrt{\frac{m^2}{2(k^2 - 1)}} (x - kt + c_0) \right], \quad k^2 - 1 > 0, \quad n < 0;$$

$$u_2(x, t) = \pm m \sqrt{\frac{2}{-n}} \operatorname{sech} \left[\sqrt{\frac{m^2}{1 - k^2}} (x - kt + c_0) \right], \quad k^2 - 1 < 0, \quad n < 0;$$

$$u_3(x, t) = \pm \frac{m}{\sqrt{-n}} \exp \left\{ \pm i \arcsin \operatorname{sech} \left[m \sqrt{\frac{2}{k^2 - 1}} (x - kt + c_0) \right] \right\},$$

$$k^2 - 1 > 0, \quad n < 0.$$

3. 由 (I), (II), (III) 获得 Duffing 方程^[1,7,10] $u_{tt} + bu + cu^3 = 0$ 的三组行波解

$$u_1(x, t) = \pm \sqrt{\frac{-b}{c}} \tanh \left[\sqrt{\frac{b}{2k^2}} (-kt + c_0) \right], \quad b > 0, \quad c < 0;$$

$$u_2(x, t) = \pm \sqrt{\frac{-2b}{c}} \operatorname{sech} \left[\sqrt{\frac{-b}{k^2}} (-kt + c_0) \right], \quad b < 0, \quad c > 0;$$

$$u_3(x, t) = \pm \sqrt{\frac{-b}{c}} \exp \left\{ \pm i \arcsin \operatorname{sech} \left[\sqrt{\frac{2b}{k^2}} (-kt + c_0) \right] \right\}, \quad b > 0, \quad c < 0.$$

4. 由 (I), (II), (III) 获得 Landau-Ginzburg-Higgs 方程^[2,7,10] $u_{tt} - u_{xx} - m^2 u + n^2 u^3 = 0$ 的三组行波解

$$u_1(x, t) = \pm \left| \frac{m}{n} \right| \tanh \left[\sqrt{\frac{m^2}{2(1 - k^2)}} (x - kt + c_0) \right], \quad k^2 - 1 < 0;$$

$$u_2(x, t) = \pm \sqrt{2} \left| \frac{m}{n} \right| \operatorname{sech} \left[\sqrt{\frac{m^2}{k^2 - 1}} (x - kt + c_0) \right], \quad k^2 - 1 > 0;$$

$$u_3(x, t) = \pm \left| \frac{m}{n} \right| \exp \left\{ \pm i \arcsin \operatorname{sech} \left[m \sqrt{\frac{2}{1 - k^2}} (x - kt + c_0) \right] \right\}, \quad k^2 - 1 < 0.$$

5. 由 (I), (II), (III) 获得 sin-Gordon 方程^[2,7,9] $u_{tt} - u_{xx} + u - \frac{1}{6} u^3 = 0$ 的三组行波解

$$u_1(x, t) = \pm \sqrt{6} \tanh \left[\sqrt{\frac{1}{2(k^2 - 1)}} (x - kt + c_0) \right], \quad k^2 - 1 > 0;$$

$$u_2(x, t) = \pm 2 \sqrt{3} \operatorname{sech} \left[\sqrt{\frac{1}{1 - k^2}} (x - kt + c_0) \right], \quad k^2 - 1 < 0;$$

$$u_3(x, t) = \pm \sqrt{6} \exp \left\{ \pm i \arcsin \operatorname{sech} \left[\sqrt{\frac{2}{k^2 - 1}} (x - kt + c_0) \right] \right\}, \quad k^2 - 1 > 0.$$

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NEW EXPLICIT AND TRAVELLING WAVE SOLUTIONS FOR A CLASS OF NONLINEAR EVOLUTION EQUATIONS*

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ABSTRACT

In this paper, with the help of Mathematica, three travelling wave solutions for a class of non-linear evolution equations $u_t + au_{xx} + bu + cu^3 = 0$ are obtained by trigonometric function method and Wu-elimination method which include new travelling wave solutions, bell soliton solutions and kink soliton solutions. Some equations such as Duffing equation, Klein-Gordon equation, Landau-Ginburg-Higgs equation and ϕ^4 equation are particular cases of the evolution equations. The method also can be applied to other nonlinear equations.

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