

# 一般格子方程新的无穷序列精确解\*

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为了获得非线性差分微分方程的无穷序列精确解, 引入一种双曲函数型辅助方程, 并给出该方程解的 Bäcklund 变换和解的非线性叠加公式. 在此基础上, 利用辅助方程与函数变换相结合的方法, 借助符号计算系统 Mathematica, 用一般格子方程为应用实例, 获得了无穷序列精确解.

**关键词:** 非线性差分微分方程, 函数变换, 双曲函数型辅助方程, 无穷序列精确解

**PACC:** 0230, 0340, 0290

## 1. 引 言

利用齐次平衡法<sup>[1]</sup>和双曲正切函数展开法<sup>[2]</sup>, 可构造非线性发展方程的精确解. 在文献<sup>[3]</sup>中把双曲正切函数展开法中的公式  $u(x, t) = U(\xi) = \sum_{j=0}^m a_j \tanh^j(\xi)$  中的  $\tanh(\xi)$  替换为 Riccati 方程  $z'(\xi) = R + z^2(\xi)$  的解  $z(\xi)$ , 并获得了非线性发展方程的新精确解. 在此基础上, 构造非线性发展方程求解领域中提出各种辅助方程法, 并获得了诸多新成果<sup>[4-27]</sup>. 但是, 这些成果主要获得了连续非线性发展方程的有限多个精确解. 随着计算机技术和非线性科学的不断发展, 在许多领域中构造了非线性离散的数学模型, 比如: 经济学领域、人口领域、生物化学领域、海洋学领域和医学领域等. 非线性离散科学的研究引起人们的关注. 最近, 在构造非线性差分微分方程精确解方面也提出试探函数法等一些求解方法<sup>[18-27]</sup>, 并获得了有限多个精确解. 理论上, “非线性发展方程存在无穷多个解”. 为了获得非线性差分微分方程的无穷序列精确解, 本文引入一种双曲函数型辅助方程, 并给出该方程解的 Bäcklund 变换和解的非线性叠加公式. 在此基础上, 利用辅助方程与函数变换相结合的方法, 借助符号计算系统 Mathematica, 用一般格子方程<sup>[18]</sup>(1)为应用实例, 获得了无穷序列精确解. 其中包括无穷序

列双曲函数解、无穷序列三角函数解、双曲函数与有理函数相结合的无穷序列解和三角函数与有理函数相结合的无穷序列解.

$$\dot{u}_n(t) = (\alpha + \beta u_n + \gamma u_n^2)(u_{n-1} - u_{n+1}). \quad (1)$$

方程(1)包含下列格子方程:

当  $\alpha = 1$  时, 方程(1)转化为 Hybrid 格子方程<sup>[19-27]</sup>.

$$\dot{u}_n(t) = (1 + \beta u_n + \gamma u_n^2)(u_{n-1} - u_{n+1}). \quad (2)$$

当  $\beta = 0, \gamma = -1$  时, 方程(1)转化为离散的 mKdV 格子方程<sup>[21-23, 27]</sup>.

$$\dot{u}_n(t) = (\alpha - u_n^2)(u_{n-1} - u_{n+1}). \quad (3)$$

当  $\alpha = \beta = 0, \gamma = 1$  时, 方程(1)转化为修正的 Volterra 格子方程<sup>[24-26]</sup>.

$$\dot{u}_n(t) = u_n^2(u_{n-1} - u_{n+1}). \quad (4)$$

当  $\alpha = 0, \beta = 1$  时, 方程(1)转化为下列修正的 Volterra 格子方程<sup>[24-26]</sup>.

$$\dot{u}_n(t) = u_n(1 + \gamma u_n)(u_{n+1} - u_{n-1}). \quad (5)$$

一般格子方程(1)包含了混合格子方程(2)等非线性差分微分方程, 获得该方程的无穷序列精确解具有重要意义.

## 2. 双曲函数型辅助方程以及解的 Bäcklund 变换和解的非线性叠加公式

为了获得非线性差分微分方程的无穷序列精

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确解,引入下列双曲函数型辅助方程以及其解的 Bäcklund 变换和解的非线性叠加公式.

$$\frac{dz(\xi)}{d\xi} = a + b \cosh(z(\xi)) + c \sinh(z(\xi)). \quad (6)$$

通过下列变换(7),把方程(6)转化为 Riccati 方程(8).

$$\begin{aligned} \sinh(z(\xi)) &= (\exp(z(\xi)) - \exp(-z(\xi)))/2, \\ \cosh(z(\xi)) &= (\exp(z(\xi)) + \exp(-z(\xi)))/2, \\ \exp(z(\xi)) &= \nu(\xi). \end{aligned} \quad (7)$$

$$\begin{aligned} \frac{d\nu(\xi)}{d\xi} &= \frac{1}{2}(b+c)\nu^2(\xi) \\ &+ a\nu(\xi) + \frac{1}{2}(b-c). \end{aligned} \quad (8)$$

对方程(8)再做下列两种变换,可以获得 Riccati 方程(11).

$$\nu(\xi) = \frac{2}{b+c} y(\xi). \quad (9)$$

$$y(\xi) = -\frac{a}{2} + \phi(\xi). \quad (10)$$

$$\frac{d\phi(\xi)}{d\xi} = \phi^2(\xi) + \frac{1}{4}(b^2 - a^2 - c^2). \quad (11)$$

## 2.1. Riccati 方程的解

经过计算获得了 Riccati 方程(11)的下列解.

$$\phi_0(\xi) = -\sqrt{-R} \tanh(\sqrt{-R}\xi) \quad (R < 0), \quad (12)$$

$$\phi_0(\xi) = -\sqrt{-R} \coth(\sqrt{-R}\xi) \quad (R < 0), \quad (13)$$

$$\phi_0(\xi) = \sqrt{R} \tan(\sqrt{R}\xi) \quad (R > 0), \quad (14)$$

$$\phi_0(\xi) = -\sqrt{R} \cot(\sqrt{R}\xi) \quad (R > 0), \quad (15)$$

$$\phi_0(\xi) = -\frac{1}{\xi} \quad (R = 0), \quad (16)$$

$$\phi_1(\xi) = \frac{BR + A\sqrt{-R} \tanh(\sqrt{-R}\xi)}{-A + B\sqrt{-R} \tanh(\sqrt{-R}\xi)} \quad (R < 0), \quad (17)$$

$$\phi_1(\xi) = \frac{-(r\sqrt{R} + CR)\cos(\sqrt{R}\xi) + (r - C\sqrt{R})\sqrt{R}\sin(\sqrt{R}\xi)}{(r - C\sqrt{R})\cos(\sqrt{R}\xi) + (r + C\sqrt{R})\sin(\sqrt{R}\xi)} \quad (R > 0), \quad (18)$$

$$\phi_1(\xi) = \frac{-3BR + 4A\sqrt{R} - 5BR\sin(2\sqrt{R}\xi) - 5A\sqrt{R}\cos(2\sqrt{R}\xi)}{3A + 4B\sqrt{R} + 5A\sin(2\sqrt{R}\xi) - 5B\sqrt{R}\cos(2\sqrt{R}\xi)} \quad (R > 0), \quad (19)$$

$$\phi_1(\xi) = \frac{-BR + A\sqrt{R}[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]}{A + B\sqrt{R}[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]} \quad (R > 0), \quad (20)$$

$$\phi_1(\xi) = \frac{\sqrt{R}[\cos(\sqrt{R}\xi) + \sin(\sqrt{R}\xi)]}{\cos(\sqrt{R}\xi) - \sin(\sqrt{R}\xi)} \quad (R > 0), \quad (21)$$

$$\phi_1(\xi) = \frac{\sqrt{R}[-2AB\sqrt{R} + (A^2 - B^2R)[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]]}{A^2 - B^2R + 2AB\sqrt{R}[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]} \quad (R > 0), \quad (22)$$

其中  $r, A, B, C, R = (b^2 - a^2 - c^2)/4, a, b, c$  是常数.

## 2.2. Riccati 方程解的 Bäcklund 变换

若  $\phi(\xi)$  是 Riccati 方程(11)的解,则下面给出的  $\bar{\phi}(\xi)$  也是 Riccati 方程(11)的解.

$$\bar{\phi}(\xi) = \frac{p + q\phi(\xi) + m\phi^2(\xi) + n\phi^3(\xi) + r\phi'(\xi) + l[\phi'(\xi)]^2}{H + L\phi(\xi) + D\phi^2(\xi) + F\phi^3(\xi) + C\phi'(\xi) + K[\phi'(\xi)]^2}, \quad (23)$$

$$\bar{\phi}(\xi) = \frac{-BR + A\phi(\xi)}{A + B\phi(\xi)}. \quad (24)$$

Riccati 方程(11)的任意解与解的 Bäcklund 变换(23), (24)相结合,迭代运用可得到 Riccati 方程(11)的无穷序列精确解.下面列出解的两种非线性叠加公式.

$$\begin{aligned}\phi_k(\xi) &= \frac{p + q\phi_{k-1}(\xi) + m\phi_{k-1}^2(\xi) + n\phi_{k-1}^3(\xi) + r\phi'_{k-1}(\xi) + l[\phi'_{k-1}(\xi)]^2}{H + L\phi_{k-1}(\xi) + D\phi_{k-1}^2(\xi) + F\phi_{k-1}^3(\xi) + C\phi'_{k-1}(\xi) + K[\phi'_{k-1}(\xi)]^2}, \\ \phi_0(\xi) &= -\sqrt{-R}\tanh(\sqrt{-R}\xi), \\ R &= \frac{1}{4}(b^2 - a^2 - c^2) \quad (b^2 - a^2 - c^2 < 0; k = 1, 2, \dots).\end{aligned}\quad (25)$$

$$\begin{aligned}\phi_k(\xi) &= \frac{p + q\phi_{k-1}(\xi) + m\phi_{k-1}^2(\xi) + n\phi_{k-1}^3(\xi) + r\phi'_{k-1}(\xi) + l[\phi'_{k-1}(\xi)]^2}{H + L\phi_{k-1}(\xi) + D\phi_{k-1}^2(\xi) + F\phi_{k-1}^3(\xi) + C\phi'_{k-1}(\xi) + K[\phi'_{k-1}(\xi)]^2}, \\ \phi_0(\xi) &= \sqrt{R}\tan(\sqrt{R}\xi), \\ R &= \frac{1}{4}(b^2 - a^2 - c^2) \quad (b^2 - a^2 - c^2 > 0; k = 1, 2, \dots).\end{aligned}\quad (26)$$

其中  $q = \frac{1}{Kl}[Ll^2 - (l^2 + K^2R)[m + r + (F + l)R]]$ ,  $n = \frac{1}{K}(Fl - l^2 - K^2R)$ ,  $H = \frac{1}{K}[Ll - l^2R - l(m + r + FR) - KR(C + KR)]$ ,  $p = R(-L + m + FR)$ ,  $D = -C + \frac{1}{K}(Fl - l^2) + \frac{1}{l}(m + r + FR)K - KR$ ,  $R = \frac{1}{4}(b^2 - a^2 - c^2)$ ;  $L, A, B, C, K, F, l, m, r, a, b, c$  是常数.

### 2.3. Riccati 方程解的非线性叠加公式

若  $\phi_1(\xi), \phi_2(\xi), \phi_3(\xi)$  是 Riccati 方程(11)的 3 个解, 则下面给出的  $\bar{\phi}(\xi)$  也是 Riccati 方程(11)的解.

$$\bar{\phi}(\xi) = \frac{R[-r\phi_1(\xi) + (p+r)\phi_2(\xi) - p\phi_3(\xi)]}{-r\phi_2(\xi)\phi_3(\xi) + \phi_1(\xi)[-p\phi_2(\xi) + (p+r)\phi_3(\xi)]}, \quad (27)$$

$$\bar{\phi}(\xi) = \frac{r\phi_2(\xi)\phi_3(\xi) - \phi_1(\xi)[P\phi_2(\xi) + (-P+r)\phi_3(\xi)]}{-r\phi_1(\xi) + (-P+r)\phi_2(\xi) + P\phi_3(\xi)}, \quad (28)$$

$$\bar{\phi}(\xi) = \frac{R[N\phi_3(\xi) - [-G + m\phi_3(\xi)]\phi_2(\xi) - \Psi(\xi)\phi_1(\xi)]}{nR\phi_3(\xi) + \Phi(\xi)\phi_1(\xi) - [(m+n)R + (G+N)\phi_3(\xi)]\phi_2(\xi)}, \quad (29)$$

这里  $\Psi(\xi) = G + N + n\phi_2(\xi) - (m+n)\phi_3(\xi)$ ,  $\Phi(\xi) = mR + N\phi_2(\xi) + G\phi_3(\xi)$ ;  $P, G, N, p, r, m, n, R = \frac{1}{4}(b^2 - a^2 - c^2)$ ,  $a, b, c$  都是常数.

Riccati 方程(11)不同 3 个解与解的非线性叠加公式(27)–(29)式相结合, 获得 Riccati 方程(11)的无穷序列复合型精确解, 这里列出解的几种非线性叠加公式.

$$\begin{aligned}\phi_k(\xi) &= \frac{R[-r\phi_{k-3}(\xi) + (p+r)\phi_{k-2}(\xi) - p\phi_{k-1}(\xi)]}{-r\phi_{k-2}(\xi)\phi_{k-1}(\xi) + \phi_{k-3}(\xi)[-p\phi_{k-2}(\xi) + (p+r)\phi_{k-1}(\xi)]} \quad (k = 3, 4, \dots), \\ \phi_0(\xi) &= \frac{-BR + A\sqrt{R}[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]}{A + B\sqrt{R}[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]},\end{aligned}\quad (30)$$

$$\phi_1(\xi) = \frac{\sqrt{R}[\cos(\sqrt{R}\xi) + \sin(\sqrt{R}\xi)]}{\cos(\sqrt{R}\xi) - \sin(\sqrt{R}\xi)},$$

$$\phi_2(\xi) = -\frac{1}{\xi}.$$

当  $k = 3$  时, 从非线性叠加公式(30), 得到 Riccati 方程(11)的如下复合型新解.

$$\phi_{31}(\xi) = \frac{\sqrt{R}[T_1(\xi) + T_2(\xi)\cos(2\sqrt{R}\xi) + T_3(\xi)\sin(2\sqrt{R}\xi)]}{T_4(\xi) + T_5(\xi)\cos(2\sqrt{R}\xi) + T_6(\xi)\sin(2\sqrt{R}\xi)}, \quad (31)$$

$$\phi_k(\xi) = \frac{R[-r\phi_{k-3}(\xi) + (p+r)\phi_{k-2}(\xi) - p\phi_{k-1}(\xi)]}{-r\phi_{k-2}(\xi)\phi_{k-1}(\xi) + \phi_{k-3}(\xi)[-p\phi_{k-2}(\xi) + (p+r)\phi_{k-1}(\xi)]} \quad (k = 3, 4, \dots),$$

$$\begin{aligned}\phi_0(\xi) &= -\sqrt{-R} \tanh(\sqrt{-R}\xi), \phi_1(\xi) = \frac{BR + A\sqrt{-R} \tanh(\sqrt{-R}\xi)}{-A + B\sqrt{-R} \tanh(\sqrt{-R}\xi)}, \\ \phi_2(\xi) &= -\frac{1}{\xi}.\end{aligned}\quad (32)$$

当  $k = 3$  时, 从非线性叠加公式(32)得到 Riccati 方程(11)的如下复合型新解.

$$\phi_{32}(\xi) = -\frac{R[-A + \sqrt{-R}(B + A\xi) \tanh(\sqrt{-R}\xi) + BR\xi \tanh^2(\sqrt{-R}\xi)]}{BR - \sqrt{-R}(-A + BR\xi) \tanh(\sqrt{-R}\xi) + AR\xi \tanh^2(\sqrt{-R}\xi)}, \quad (33)$$

$$\phi_k(\xi) = \frac{R[-r\phi_{k-3}(\xi) + (p+r)\phi_{k-2}(\xi) - p\phi_{k-1}(\xi)]}{-r\phi_{k-2}(\xi)\phi_{k-1}(\xi) + \phi_{k-3}(\xi)[-p\phi_{k-2}(\xi) + (p+r)\phi_{k-1}(\xi)]} \quad (k = 3, 4, \dots),$$

$$\phi_0(\xi) = \frac{\sqrt{R}[\cos(\sqrt{R}\xi) + \sin(\sqrt{R}\xi)]}{\cos(\sqrt{R}\xi) - \sin(\sqrt{R}\xi)},$$

$$\phi_1(\xi) = \frac{\sqrt{R}[-2AB\sqrt{R} + (A^2 - B^2R)[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]]}{A^2 - B^2R + 2AB\sqrt{R}[\sec(2\sqrt{R}\xi) + \tan(2\sqrt{R}\xi)]},$$

$$\phi_2(\xi) = -\frac{1}{\xi}.\quad (34)$$

当  $k = 3$  时, 从非线性叠加公式(34), 得到 Riccati 方程(11)的如下复合型新解.

$$\phi_{33}(\xi) = -\frac{\sqrt{R}[\varphi_1(\xi) + \varphi_2(\xi)\cos(2\sqrt{R}\xi) + \varphi_3(\xi)\sin(2\sqrt{R}\xi)]}{\varphi_4(\xi) - \varphi_3(\xi)\cos(2\sqrt{R}\xi) + \varphi_2(\xi)\sin(2\sqrt{R}\xi)}, \quad (35)$$

其中  $T_5(\xi) = (-A + BR\xi)$ ,  $T_6(\xi) = -\sqrt{R}(B + A\xi)$ ,  $\varphi_1(\xi) = -A^2 + B^2R - 2ABR\xi$ ,  $T_1(\xi) = (A + BR\xi)$ ,  $T_2(\xi) = (B + A\xi)$ ,  $T_3(\xi) = \sqrt{R}(-A + BR\xi)$ ,  $T_4(\xi) = \sqrt{R}(B - A\xi)$ ,  $\varphi_2(\xi) = \sqrt{R}(-2AB - A^2\xi + B^2R\xi)$ ,  $\varphi_3(\xi) = A^2 - B^2R - 2ABR\xi$ ,  $\varphi_4(\xi) = \sqrt{R}(2AB - A^2\xi + B^2R\xi)$ . 我们迭代运用 Riccati 方程(11)的 Bäcklund 变换(23), (24)和解的非线性叠加公式(27)——(29)和变换(7), (9), (10)获得双曲函数型辅助方程(6)无穷序列双曲函数解、无穷序列三角函数解和无穷序列复合型精确解(这里只列出解的几种非线性叠加公式).

### 3. 一般格子方程的无穷序列精确解

下面用双曲函数型辅助方程与函数变换相结合的方法, 构造一般格子方程的无穷序列精确解.

将  $u_n(n, t) = u_n(\xi)$ ,  $u_{n+1}(n, t) = u_{n+1}(\xi)$ ,  $u_{n-1}(n, t) = u_{n-1}(\xi)$ ,  $\xi = dn + \omega t$  代入(1)后得到下列方程

$$\omega u'_n(\xi) = (\alpha + \beta u_n(\xi))$$

$$+ \gamma u_n^2(\xi))(u_{n-1}(\xi) - u_{n+1}(\xi)). \quad (36)$$

方程(36)的解取为下列形式.

$$\begin{aligned}u_n(\xi) &= g_0 + g_1 \sinh(z(\xi)) \\ &\quad + f_1 \cosh(z(\xi)),\end{aligned}\quad (37)$$

$$\begin{aligned}u_{n+1}(\xi) &= g_0 + g_1(\sinh(z(\xi)) \cosh(d) \\ &\quad - \sinh(d) \cosh(z(\xi)) + d + 1) \\ &\quad + f_1(\sinh(z(\xi)) \cosh(d) \\ &\quad + \sinh(d) \cosh(z(\xi)) \\ &\quad + d + 1),\end{aligned}\quad (38)$$

$$\begin{aligned}u_{n-1}(\xi) &= g_0 + g_1(\sinh(z(\xi)) \cosh(d) \\ &\quad + \sinh(d) \cosh(z(\xi)) + d) \\ &\quad + f_1(\sinh(z(\xi)) \cosh(d) \\ &\quad - \sinh(d) \cosh(z(\xi)) + d).\end{aligned}\quad (39)$$

将(6), (37), (38), (39)一起代入(36), 并令  $\cosh^i(z(\xi)) \sinh^j(z(\xi))$ , ( $i = 0, 1; j = 0, 1, 2, 3$ ) 的系数为零后得到一个非线性代数方程组(限于篇幅未列出). 用符号计算系统 Mathematica 求出该方程组的如下解:

$$a = \mp \frac{2f_1 \sqrt{\beta^2 + 4\gamma(-\alpha + \gamma f_1^2)}}{\omega},$$

$$b = -\frac{4\gamma f_1^2}{\omega},$$

$$\begin{aligned}
c &= 0, \\
g_0 &= \frac{-\beta \pm \sqrt{\beta^2 + 4\gamma(-\alpha + \gamma f_1^2)}}{2\gamma}, \\
g_1 &= f_1; \\
a &= -\frac{cf_1(\beta + 2\gamma g_0)}{\alpha + \beta g_0 + \gamma(-f_1^2 + g_0^2)}, \\
b &= -\frac{c[\alpha + \beta g_0 + \gamma(f_1^2 + g_0^2)]}{\alpha + \beta g_0 + \gamma(-f_1^2 + g_0^2)}, \\
\omega &= \frac{2[\alpha + \beta g_0 + \gamma(-f_1^2 + g_0^2)]}{c}, \\
g_1 &= f_1.
\end{aligned} \tag{40}$$

将(40),(41)分别与(7)一起代入(37)即可得到一般格子方程的下列形式的解:

$$\begin{aligned}
u_n(n, t) = u_n(\xi) &= \frac{-\beta \pm \sqrt{\beta^2 + 4\gamma(-\alpha + \gamma f_1^2)}}{2\gamma} \\
&+ f_1 \nu(dn + \omega t);
\end{aligned} \tag{42}$$

$$\begin{aligned}
u_n(n, t) = u_n(\xi) &= g_0 + f_1 \nu \left( dn \right. \\
&\left. + \frac{2[\alpha + \beta g_0 + \gamma(-f_1^2 + g_0^2)]}{c} t \right).
\end{aligned} \tag{43}$$

利用下列解的非线性叠加公式,可以获得一般格子方程的无穷序列精确解:

$$\begin{aligned}
\nu_k(\xi) &= \frac{2}{b+c} \left[ \phi_k(\xi) - \frac{a}{2} \right], \\
\phi_k(\xi) &= \frac{p + q\phi_{k-1}(\xi) + m\phi_{k-1}^2(\xi) + n\phi_{k-1}^3(\xi) + r\phi'_{k-1}(\xi) + l[\phi'_{k-1}(\xi)]^2}{H + L\phi_{k-1}(\xi) + D\phi_{k-1}^2(\xi) + F\phi_{k-1}^3(\xi) + C\phi'_{k-1}(\xi) + K[\phi'_{k-1}(\xi)]^2}, \\
\phi_0(\xi) &= -\sqrt{-R} \tanh(\sqrt{-R}\xi),
\end{aligned} \tag{44}$$

$$R = \frac{1}{4}(b^2 - a^2 - c^2) \quad (b^2 - a^2 - c^2 < 0; k = 1, 2, \dots).$$

$$\begin{aligned}
\nu_k(\xi) &= \frac{2}{b+c} \left[ \phi_k(\xi) - \frac{a}{2} \right], \\
\phi_k(\xi) &= \frac{p + q\phi_{k-1}(\xi) + m\phi_{k-1}^2(\xi) + n\phi_{k-1}^3(\xi) + r\phi'_{k-1}(\xi) + l[\phi'_{k-1}(\xi)]^2}{H + L\phi_{k-1}(\xi) + D\phi_{k-1}^2(\xi) + F\phi_{k-1}^3(\xi) + C\phi'_{k-1}(\xi) + K[\phi'_{k-1}(\xi)]^2}, \\
\phi_0(\xi) &= \sqrt{R} \tan(\sqrt{R}\xi),
\end{aligned} \tag{45}$$

$$R = \frac{1}{4}(b^2 - a^2 - c^2) \quad (b^2 - a^2 - c^2 > 0; k = 1, 2, \dots).$$

$$\begin{aligned}
\nu_k(\xi) &= \frac{2}{b+c} \left[ \phi_k(\xi) - \frac{a}{2} \right], \\
R &= \frac{1}{4}(b^2 - a^2 - c^2) \quad (b^2 - a^2 - c^2 > 0; k = 1, 2, \dots), \\
\phi_k(\xi) &= \frac{p + q\phi_{k-1}(\xi) + m\phi_{k-1}^2(\xi) + n\phi_{k-1}^3(\xi) + r\phi'_{k-1}(\xi) + l[\phi'_{k-1}(\xi)]^2}{H + L\phi_{k-1}(\xi) + D\phi_{k-1}^2(\xi) + F\phi_{k-1}^3(\xi) + C\phi'_{k-1}(\xi) + K[\phi'_{k-1}(\xi)]^2}, \\
\phi_0(\xi) &= \frac{\sqrt{R}[T_1(\xi) + T_2(\xi)\cos(2\sqrt{R}\xi) + T_3(\xi)\sin(2\sqrt{R}\xi)]}{T_4(\xi) + T_5(\xi)\cos(2\sqrt{R}\xi) + T_6(\xi)\sin(2\sqrt{R}\xi)}.
\end{aligned} \tag{46}$$

$$\begin{aligned}
\nu_k(\xi) &= \frac{2}{b+c} \left[ \phi_k(\xi) - \frac{a}{2} \right], \\
R &= \frac{1}{4}(b^2 - a^2 - c^2) \quad (b^2 - a^2 - c^2 < 0; k = 1, 2, \dots), \\
\phi_k(\xi) &= \frac{p + q\phi_{k-1}(\xi) + m\phi_{k-1}^2(\xi) + n\phi_{k-1}^3(\xi) + r\phi'_{k-1}(\xi) + l[\phi'_{k-1}(\xi)]^2}{H + L\phi_{k-1}(\xi) + D\phi_{k-1}^2(\xi) + F\phi_{k-1}^3(\xi) + C\phi'_{k-1}(\xi) + K[\phi'_{k-1}(\xi)]^2}, \\
\phi_0(\xi) &= -\frac{R[-A + \sqrt{-R}(B + A\xi)]\tanh(\sqrt{-R}\xi) + BR\xi\tanh^2(\sqrt{-R}\xi)}{BR - \sqrt{-R}(-A + BR\xi)\tanh(\sqrt{-R}\xi) + AR\xi\tanh^2(\sqrt{-R}\xi)}.
\end{aligned} \tag{47}$$

其中  $q = \frac{1}{Kl}[Ll^2 - (l^2 + K^2R)[m + r + (F + l)R]]$ ,  
 $n = \frac{1}{K}(Fl - l^2 - K^2R)$ ,  $H = \frac{1}{K}[Ll - l^2R - l(m + r + FR) - KR(C + KR)]$ ,  $p = R(-L + m + FR)$ ,  $D = -C + \frac{1}{K}(Fl - l^2) + \frac{1}{l}(m + r + FR)K - KR$ ,  $R = \frac{1}{4}(b^2 - a^2 - c^2)$ ,  $T_1(\xi) = (A + BR\xi)$ ,  $T_2(\xi) = (B + A\xi)$ ,  $T_3(\xi) = \sqrt{R}(-A + BR\xi)$ ,  $T_4(\xi) = \sqrt{R}(B - A\xi)$ ,  $T_5(\xi) = (-A + BR\xi)$ ,  $T_6(\xi) = -\sqrt{R}(B + A\xi)$ ,  $L, A, B, C, K, F, l, m, r, a, b, c$  是任意常数.

将解的非线性叠加公式(44)–(47)式来确定的无穷序列解, 分别代入(42), (43)式后得到一般格子方程的无穷序列双曲函数解、无穷序列三角函数解、双曲函数与有理函数相结合的无穷序列解和三角函数与有理函数相结合的无穷序列解.

## 4. 结 论

非线性发展方程求解方法的研究, 是孤立子理论的重要研究课题之一. 数学物理学家们, 为了获得非线性发展方程的新解而提出许多新方法, 并获得诸多新成果. 这些方法当中辅助方程法值得人们的关注. 以往辅助方程法有关的文献只获得了非线性发展方程的有限多个精确解. 没有得到无穷序列精确解. 比如: 文献[18–27]获得了 Hybrid 格子方程等非线性差分微分方程的双曲函数和三角函数形式的有限多个精确解. 本文为了获得非线性差分微分方程的无穷序列精确解, 给出一种双曲函数型辅助方程, 并获得了该方程解的 Bäcklund 变换和解的非线性叠加公式, 利用辅助方程与函数变换相结合的方法, 借助符号计算系统 Mathematica, 构造了一般格子方程(1)的四种类型的无穷序列新精确解.

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# New infinite sequence exact solutions to the general lattice<sup>\*</sup>

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## Abstract

To seek infinite sequence exact solutions to nonlinear difference-differential equations, a kind of auxiliary equation of hyperbolic function type is introduced, and the Bäcklund transformation of the solutions to the equation and the formula of nonlinear superposition of solutions are presented. Based on these and combining auxiliary equation with function transformation, the general lattice equation is chosen as example for which the infinite sequence exact solutions to the equation are obtained.

**Keywords:** nonlinear difference-differential equation, function transformation, auxiliary equation of hyperbolic function type, infinite sequence exact solution

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