

A VERIFICATION OF ANDRADE'S FORMULA OF VISCOSITY

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The well known formula of Andrade¹ for the viscosity η of a liquid at its freezing point is

$$\eta = \frac{4}{3} \nu m / \sigma, \quad (1)$$

where m is the molecular mass, σ the average intermolecular distance (center-to-center), and ν the frequency of molecular vibration of the liquid. Denoting molar volume by V , Avogadro number by N , freezing point by T , gas constant by R , molecular weight by M , and molecular radius by a , we may put

$$\sigma = (V/N)^{1/3}$$

$$f = 1 + \frac{l}{a} (0.864 + 0.29e^{-1.25a/l}), \quad (5)$$

and, as given by Lindemann's theory² of fusion,

$$\nu = CT^{1/2} M^{-1/2} V^{-1/3},$$

where

$$C = \frac{\sqrt{2}}{\pi \rho} R^{1/2} N^{1/3}, \quad (2)$$

ρ being a fraction defined by

$$2a = \sigma(1 - \rho). \quad (3)$$

Equation (3) means that $\sigma\rho$ is the shortest face-to-face distance between the assumed spheres representing two neighboring molecules when they are at their mean positions.

In support to Andrade's theory, we have found a simple, though not very rigorous, way of verifying formula (1). The motion of each molecule may be treated as one in a discontinuous medium composed of all the rest molecules. The resistance offered by this medium to each molecule moving through it with velocity v is^{3,4}

$$R = 4\pi a \eta v / f, \quad (4)$$

where

l being the mean free path of the molecules. If we set l equal to $\sigma\rho$, i.e., twice the amplitude of vibration in Lindemann's theory of fusion, equation (3) gives

$$\frac{l}{a} = \frac{2\rho}{1 - \rho}. \quad (6)$$

Adopting Gruneisen's value⁵ 2.8×10^{12} for C , equation (2) yields $\rho = 0.124$. By (6) and (5), we have $f = 1.27$ for any liquid at its freezing point. The equation of motion for any chosen direction x is

$$m \frac{d^2x}{dt^2} = -\alpha \frac{dx}{dt} + X,$$

where X is the force causing motion and

$$\alpha = 4\pi a \eta / f. \quad (7)$$

As in Langevin's theory of Brownian movement, multiplying through by x changes this equation into

$$\begin{aligned} \frac{1}{2} m \frac{d^2(x^2)}{dt^2} - m \left(\frac{dx}{dt} \right)^2 = \\ -\frac{\alpha}{2} \frac{d(x^2)}{dt} + Xx. \end{aligned}$$

Taking the statistical average of this equation over all the molecules, we have

$$m \frac{d^2(x^2)}{dt^2} - 2kT = -\alpha \frac{d(x^2)}{dt}.$$

since $\overline{Xx} = 0$ and, by the principle of energy equipartition, $\frac{1}{2} m \left(\frac{dx}{dt} \right)^2 = \frac{1}{2} kT$, k being the Boltzmann constant. The solution to this equation is

$$\frac{d(x^2)}{dt} = \frac{2kT}{\alpha} \left(1 - e^{-\frac{\alpha}{m} t} \right).$$

If t is greater than 10^{-14} second, the exponential term in the bracket is negligible compared to unity for all molecules except hydrogen. In this case, neglect of this term gives

$$\int_0^x d(x^2) = \frac{2kT}{\alpha} \int_0^t dt,$$

$$\text{or } \frac{x}{t} = \frac{2kT}{\alpha t}.$$

If t is the time between two successive collisions, i.e., equal to $1/2v$, then x/t is the average component speed, \bar{v} , of the molecules along the x -direction chosen. I.e.,

$$\begin{aligned} \left(\frac{x}{t} \right)^2 = (\bar{v})^2 &= \frac{\left[\int_0^\infty ve^{-\frac{1}{2} \frac{mv^2}{kT}} dv \right]^2}{\left[\int_0^\infty e^{-\frac{1}{2} \frac{mv^2}{kT}} dv \right]} \\ &= \frac{2kT}{\pi m}. \end{aligned}$$

$$\therefore \pi m = \alpha t = \frac{\alpha}{2v},$$

$$\text{or } \alpha = 2\pi m v.$$

Substituting this expression of α into (7), we have

$$\eta = \frac{mvf}{2\alpha},$$

which, by (3), becomes

$$\eta = \frac{mvf}{\alpha(1-\rho)}.$$

Putting in the numerical values of f and ρ previously given, we have

$$\eta = 1.45 \frac{mv}{\sigma}. \quad (8)$$

which is only 1.09 times that of Andrade's expression (1).

For diatomic liquids, the molecular shape is not spherical and therefore

equation (8) is not expected to give in the case of hydrogen. This is in accurate values. It however yields the the present theory due to the fact that correct tendency of change for different for hydrogen molecules as shown in the following table:

$$e^{-\frac{av}{m}} = e^{-\frac{a}{2mv}} = e^{-0.79} = 0.45,$$

Molecule	Freezing point	calculated	observed
Cl ₂	171°K	7.7 c.p.	10.3 c.p. ⁶
Br ₂	265.7	9.8	13.8 ⁶
I ₂	286	14.5	23.2 ⁶
O ₂	54	3.1	8.1 ⁷
H ₂	14	0.037	0.013 ⁸

We note that the discrepancy is largest

which is certainly not negligible compared to unity so that the simple solution here used is not applicable.

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1. Andrade, E. N. da C., *Phil. Mag.* **17**, (1934) 497, 698.
2. Lindemann, *Phys. Zeits.* **11**, (1912) 609.
3. Herzog, R.O., and Kudar, H., *Zeits. f. Phys.*, **80**, (1933) 217.
4. Davies, C.N., *Proc. Roy. Soc.* **57**, (1945) 259.
5. Gruneisen, *Ann. der Phys.*, **39**, (1912) 257.
6. Steacie and Johnson, *J. Am. Chem. Soc.*, **57**, (1925) 754.
7. Rudeko and Schubnikow, *Phys. Z. Sowjetunion*, **6**, (1934) 470.
8. Verschaffelt, J.E., *Comm. Phys. Lab. Leiden*, **150**, b, (1917) 16.